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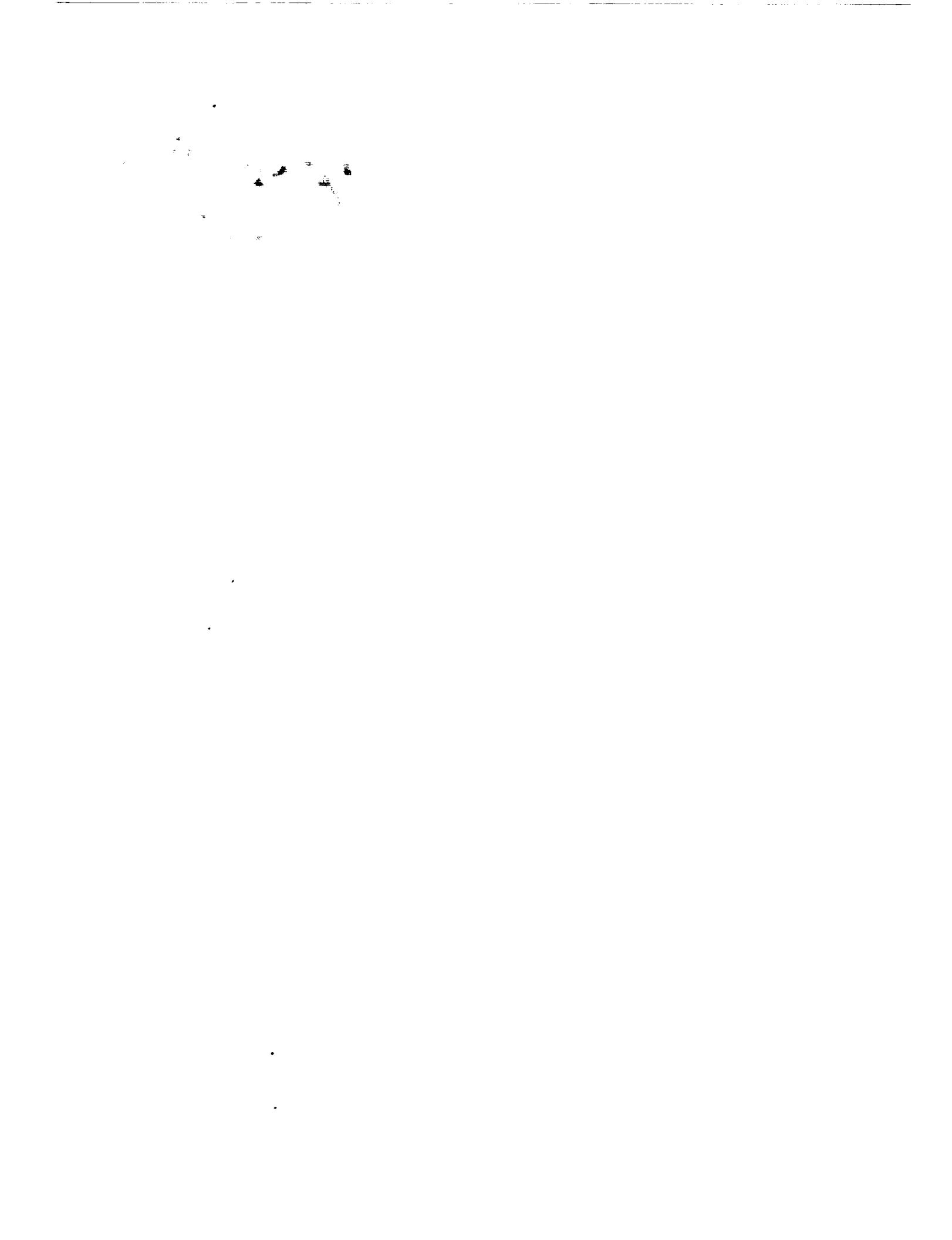
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NUMERICAL SOLUTION OF AXIALLY SYMMETRIC POISSON  
EQUATION; THEORY AND APPLICATION TO  
ION-THRUSTOR ANALYSIS

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NATIONAL AERONAUTICS AND SPACE ADMINISTRATION  
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NUMERICAL SOLUTION OF AXIALLY SYMMETRIC POISSON  
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SUMMARY

A numerical solution of the axially symmetric Poisson equation for mixed boundary conditions is presented, and the theory and application to an ion-thruster analysis is discussed.

A method of solution is developed in which the differential equation is replaced by finite difference equations, and the properties of the resulting matrix are studied. The theory of matrix "regular splittings" is applied. The Cyclic Chebyshev Semi-Iterative method is used to solve the matrix equation, and an estimate of an optimum relaxation factor is given. The Poisson equation is solved by a method of successive approximations.

The numerical method is demonstrated with an example of a hollow-cylindrical ion beam. The program for an IBM 7090 computer is also included.

INTRODUCTION

The numerical method of solution of the space-charge-flow problem or, mathematically, the solution of the two-dimensional Poisson equation is developed and presented in reference 1. Application of this two-dimensional method together with its extension involving curved ion-emitter surfaces is demonstrated in reference 2. Good agreement of this method with available experimental data is also reported in reference 2. A method of analysis of an axially symmetric ion beam in an ion thruster is discussed herein. The material presented represents a logical extension of the work described in references 1 and 2.

After boundary conditions of finite extent are stipulated, the Poisson equation in a continuous form is replaced by finite-difference equations. The finite-difference equations give rise to a real symmetric matrix by virtue of the axial symmetry of the problem. The properties and convenient partitioning of the resulting real symmetric matrix are discussed in detail. The solution of the matrix equation is accomplished by the use of the Cyclic Chebyshev Semi-Iterative method. The general procedure for obtaining a solution of the axially symmetric Poisson equation is essentially the same as that described in

reference 1. The space-charge-density-distribution function is not known a priori; and therefore, the Laplacian equation is solved to determine a first approximation of the potential distribution in the bounded region. The equations of motion are used together with this initial potential distribution, to obtain the ion trajectories and the initial values of the space-charge-density function. A series of successive approximations then gives the solution of the Poisson equation.

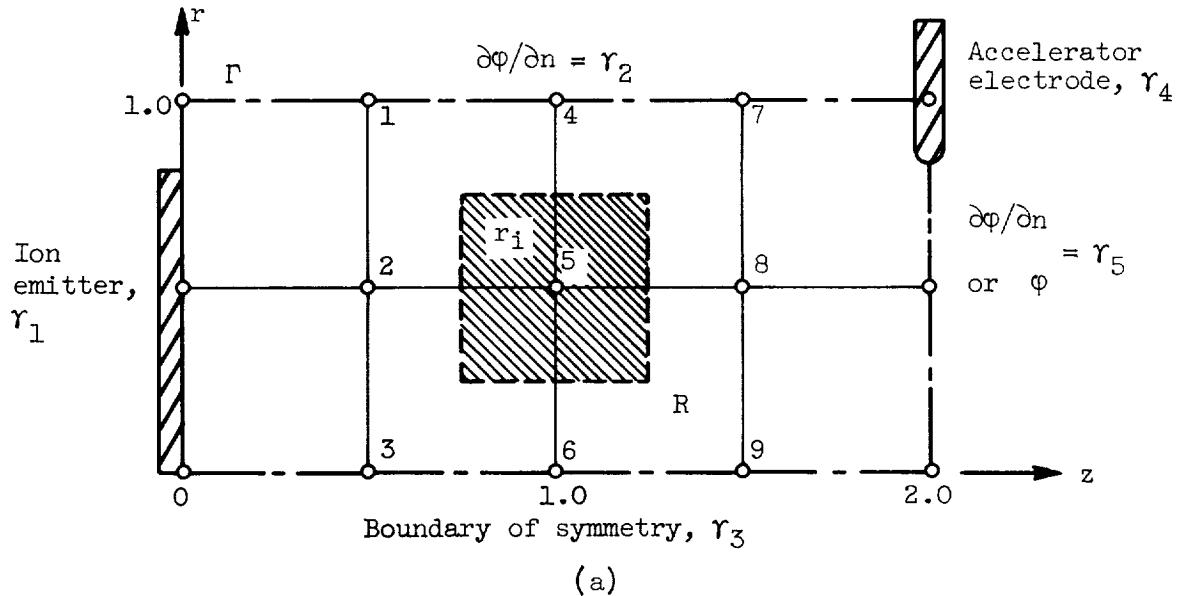
The theory of the numerical method is demonstrated in an example problem. The problem, which was analyzed, is the space-charge-limited flow in a hollow-cylindrical-beam ion thruster that has a cylindrical ion emitter. The solution was obtained with the aid of an IBM 7090 computer. The computer program is described in appendix C by Carl D. Bogart. No effort was made to optimize the computer program other than the method of obtaining rapid convergence to the solution of the matrix equation. The electrode coordinates together with the method of determination of electrode shapes are given in reference 3.

#### STATEMENT OF PROBLEM AND NUMERICAL ANALYSIS

The theory of a numerical method of solving the axially symmetric Poisson equation is presented for the steady-state space-charge-limited flow of an ion beam in an ion thruster. In this section, a mathematical model is determined from a hypothetical axially symmetric ion thruster, and the finite-difference equations that give rise to the matrix equation are established. The solution of the matrix equation is then developed.

#### Mathematical Model

A region considered for presentation of the theory of the numerical solution of the axially symmetric Poisson equation is shown in sketch (a). The external



boundary represents a hypothetical ion thruster with a disk emitter and a cylindrical ion beam extending in the z-direction. The location of the downstream boundary  $\gamma_5$  is arbitrarily shown in the plane of the accelerator. For analysis of a real thruster configuration, the location of this boundary will normally be further downstream, as is described in the example problem.

The Poisson equation in cylindrical coordinates is in the form

$$-\nabla^2 \varphi(r, \theta, z) = \frac{1}{\epsilon_0} \rho(\varphi, r, \theta, z) \quad (1)$$

The external boundary  $\Gamma$  of the region  $R$  satisfies the equation

$$\alpha\varphi(r, \theta, z) + \beta \frac{\partial\varphi(r, \theta, z)}{\partial n} = \gamma_i \quad \text{for } i=1, 2, 3, 4, 5, \alpha \neq \beta \quad (2)$$

Two sets of values of  $\alpha$  and  $\beta$  apply: 1,0 and 0,1. Values of  $\alpha = 1$  and  $\beta = 0$  correspond to Dirichlet boundary conditions, while  $\alpha = 0$  and  $\beta = 1$  correspond to Neumann boundary conditions. (All symbols are defined in appendix A.)

The potential-distribution function is  $\varphi(r, \theta, z)$  and the space-charge-density-distribution function is  $\rho(\varphi, r, \theta, z)$  both of which are continuous inside the region  $R$ . The space-charge-density distribution function  $\rho$  is nonnegative for positive ion flow and is not known a priori. It depends on the potential-distribution function  $\varphi$ , which must satisfy equation (1) and the conditions of equation (2) on the exterior boundary  $\Gamma$  of the region  $R$ .

The numerical solution of equation (1) is accomplished by overlaying a discrete number of mesh points on the region  $R$ , as shown in sketch (a). The uniformity of mesh spacing overlayed on region  $R$  is not essential but was chosen for simplification of the matrix coefficients. The size of mesh spacing is arbitrary and depends on the physical dimensions of the problem and the anticipated potential gradients in certain portions of the region  $R$ . For that reason it is possible to have several nets of uniform mesh spacing, as described in reference 1. Next, the differential equation (eq. (1)) is replaced by the finite-difference equations that satisfy the conditions of a subregion, that is, the influence area of each mesh point.

### Finite-Difference Equations - Five-

#### Point-Formula Approximation

The Poisson equation (cylindrical coordinates) for the discrete case is

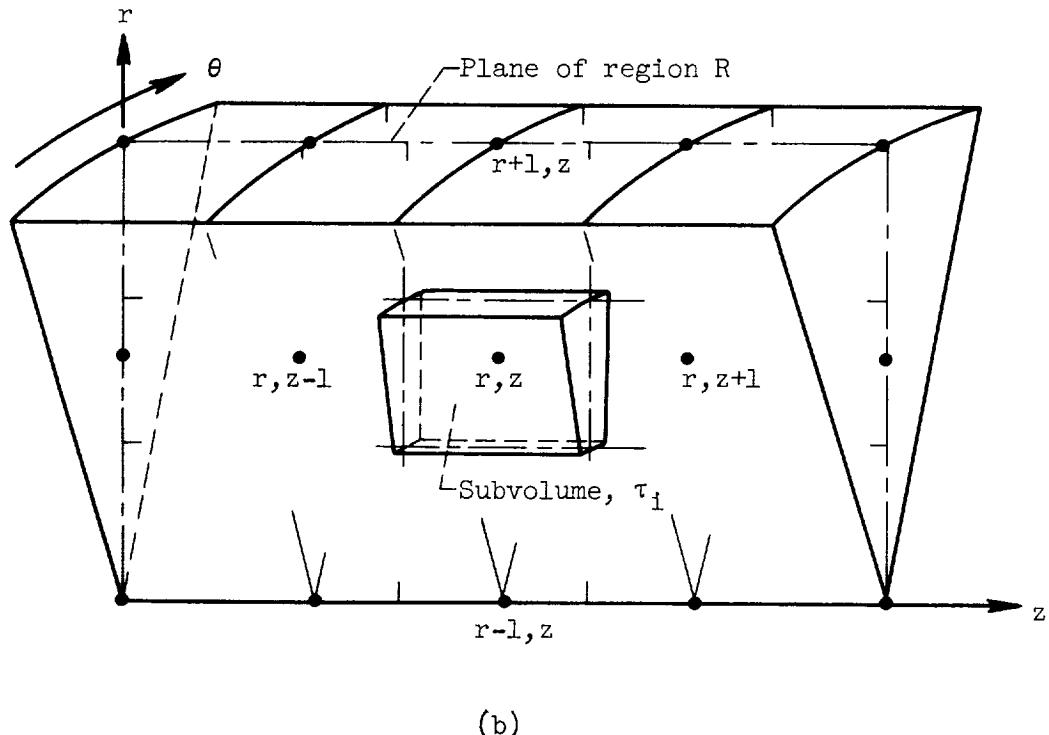
$$-\nabla^2 w(r, \theta, z) = f(w, r, \theta, z) \quad (3)$$

For each of the subvolumes  $\tau_i$  surrounding the internal point  $r, z$  (shown in sketch (b)), the numerical approximation to equation (3) is given by the five-

point-formula approximation as follows:

$$r(4w_{r,z} - w_{r,z+1} - w_{r,z-1} - w_{r+1,z} - w_{r-1,z}) + \frac{h}{2} (w_{r-1,z} - w_{r+1,z}) = rh^2 f_{r,z} \quad (4)$$

The derivation of the five-point-formula approximation from Green's three-dimensional theorem can be found in reference 4. For completeness, it is included in appendix B.



Because of the axial symmetry, equation (4) represents the finite-difference approximation of equation (3) in the  $r, z$  plane (identified as  $R$  in sketches (a) and (b)). For example, writing equation (4) for point 5 of sketch (a) gives

$$r(4w_5 - w_8 - w_2 - w_4 - w_6) + \frac{h}{2} (w_6 - w_4) = rh^2 f_5 \quad (5)$$

The five-point-formula approximation of equation (3) can be expressed for each mesh point surrounded by subregion  $r_i$  for  $i = 1, 2, \dots, N$ . The truncation error of this approximation is of the order  $h^2$ , that is,  $O(h^2)$ . For subregions along the external boundaries, the five-point-formula approximation of equation (3) must be modified to account for the particular boundary being considered. Various forms of the modified approximation formula are given in appendix B.

The finite-difference approximation of equation (3) for each subregion gives rise to a set of linear algebraic equations. For  $N$  mesh points in region  $R$  there are  $N$  linear algebraic equations with  $N$  unknowns. This set of equations can be expressed in matrix form and is discussed next.

### Matrix Equation

If the ordering of mesh points is as shown in sketch (a), the  $N$  linear algebraic equations with  $N$  unknowns can be written in matrix form as

$$Aw = k \quad (6)$$

where  $w$  is a column vector representing the discrete potential distribution  $w_{1,2}, \dots, N$ ,  $k$  is a column vector consisting of the discrete space-charge-density distribution  $f_{1,2}, \dots, N$  and, when applicable, the boundary values  $\gamma_{1,2,3,4,5}$ . The resulting matrix  $A$  is an  $N$  by  $N$  real symmetric matrix. For the example given in sketch (a), matrix  $A$  takes the form

$$A = \left[ \begin{array}{ccc|ccc|ccc} \frac{13}{8} & -\frac{3}{4} & 0 & -\frac{7}{16} & 0 & 0 & 0 & 0 & 0 \\ -\frac{3}{4} & 2 & -\frac{1}{4} & 0 & -\frac{1}{2} & 0 & 0 & 0 & 0 \\ 0 & -\frac{1}{4} & \frac{3}{8} & 0 & 0 & -\frac{1}{16} & 0 & 0 & 0 \\ \hline -\frac{7}{16} & 0 & 0 & \frac{13}{8} & -\frac{3}{4} & 0 & -\frac{7}{16} & 0 & 0 \\ 0 & -\frac{1}{2} & 0 & -\frac{3}{4} & 2 & -\frac{1}{4} & 0 & -\frac{1}{2} & 0 \\ 0 & 0 & -\frac{1}{16} & 0 & -\frac{1}{4} & \frac{3}{8} & 0 & 0 & -\frac{1}{16} \\ \hline 0 & 0 & 0 & -\frac{7}{16} & 0 & 0 & \frac{13}{8} & -\frac{3}{4} & 0 \\ 0 & 0 & 0 & 0 & -\frac{1}{2} & 0 & -\frac{3}{4} & 2 & -\frac{1}{4} \\ 0 & 0 & 0 & 0 & 0 & -\frac{1}{16} & 0 & -\frac{1}{4} & \frac{3}{8} \end{array} \right]$$

The solution of the Poisson equation (eq. (3)) now has been reduced to the numerical solution of equation (6). The diagonal entries of matrix  $A$  are positive, whereas the off-diagonal entries are nonpositive. It can be proved that the real symmetric matrix  $A$  is a Stieltjes matrix, as defined in reference 4,

and thus matrix  $A$  has an inverse  $A^{-1} > 0$  (therefore  $A$  is a nonsingular matrix), which ensures that the solution of equation (6) is unique. It should be noted that it may be necessary to multiply some of the finite-difference equations by appropriate scaling factors in order to ensure that matrix  $A$  is symmetric.

Because of the different values of the main diagonal entries, it is no longer convenient, as it was in the two-dimensional case of references 1 and 2, to premultiply the matrix  $A$  by a positive diagonal matrix  $D$  such that  $DA$  is a matrix with unity on its main diagonal. This step would destroy the symmetry of matrix  $A$  to which the following discussion is directed.

The solution of the matrix equation (eq. (6)) is based on the theory of "regular splittings" as discussed in reference 4. Using the method of this theory allows expressing matrix  $A$  as

$$A = B - C \quad (7)$$

where  $B$  and  $C$  are again  $N$  by  $N$  matrices that would be written for the example as

$$B = \begin{bmatrix} b_{1,1} & 0 & 0 \\ 0 & b_{2,2} & 0 \\ 0 & 0 & b_{3,3} \end{bmatrix}$$

where  $b_{i,i}$  represents the diagonal block matrices of the example matrix  $A$  and

$$C = \begin{bmatrix} 0 & c_{1,2} & 0 \\ c_{2,1} & 0 & c_{2,3} \\ 0 & c_{3,2} & 0 \end{bmatrix}$$

where  $c_{i,j}$  represents the off-diagonal blocks of  $A$ . Both matrices  $B^{-1}$  and  $C$  have nonnegative entries that satisfy the definition of regular splittings of matrix  $A$  according to reference 4. The theory of regular splittings is a powerful means by which a solution of matrix equations involving matrices more general than tridiagonal matrices can be accomplished very efficiently. This method could be applied as well to the solution of the matrix equation reported in references 1 and 2.

#### SOLUTION OF MATRIX EQUATION

As mentioned in the previous section, the solution of the axially symmetric

Poisson equation lies within the numerical solution of the matrix equation (eq. (6)). If equation (7) is substituted into the matrix equation (eq. (6)), it follows that

$$A\underline{w} = (B - C)\underline{w} = \underline{k}$$

or equivalently

$$\underline{w} = B^{-1}C\underline{w} + B^{-1}\underline{k} \quad (8)$$

Identifying  $D = B^{-1}C$  and  $\underline{g} = B^{-1}\underline{k}$  and substituting in equation (8) yields

$$\underline{w} = D\underline{w} + \underline{g} \quad (9)$$

Equation (9) is now used to solve the matrix equation (eq. (6)).

#### Cyclic Chebyshev Semi-Iterative Method

The solution of equation (9) is based on the fact that matrix  $A$  of equation (6) is a Stieltjes matrix, furthermore, that matrices  $B$  and  $C$  are defined as a regular splitting of  $A$ , and that  $B$  is symmetric and positive definite. The iteration matrix  $D = B^{-1}C$  of equation (9) is then irreducible, non-negative, and has real eigenvalues, since  $D$  is similar to a symmetrical matrix and is convergent, as shown in reference 4.

Now to explore the cyclic properties of the iteration matrix  $D$ . Similarly in the case of matrix  $M$  in reference 1 for the two-dimensional solution of the Poisson equation, equation (9) can be written as

$$\begin{bmatrix} \underline{w}_1 \\ \underline{w}_2 \end{bmatrix} = \begin{bmatrix} 0 & D_2 \\ D_1 & 0 \end{bmatrix} \begin{bmatrix} \underline{w}_1 \\ \underline{w}_2 \end{bmatrix} + \begin{bmatrix} \underline{g}_1 \\ \underline{g}_2 \end{bmatrix} \quad (10)$$

where now the subscript 1 is associated with odd-number lines and 2 with even-number lines. By the Cyclic Chebyshev Semi-Iterative method, the vector components of equation (10) can be written as they were in reference 1 for the two-dimensional case, namely,

$$\left. \begin{aligned} \underline{w}_1^{2m+1} &= \omega_{2m+1} \left( D_2 \underline{w}_2^{2m} + \underline{g}_1 - \underline{w}_1^{2m-1} \right) + \underline{w}_1^{2m-1} && \text{for } m \geq 1 \\ \underline{w}_2^{2m+2} &= \omega_{2m+2} \left( D_1 \underline{w}_1^{2m+1} + \underline{g}_2 - \underline{w}_2^{2m} \right) + \underline{w}_2^{2m} && \text{for } m \geq 0 \end{aligned} \right\} \quad (11)$$

where, for  $m = 0$

$$\underline{w}_1^1 = D_2 \underline{w}_2^0 + \underline{g}_1$$

Again, this iterative method requires only the initial guess of the single vector component  $\underline{w}_2^0$ , and the method of solution does not require any more computer storage space than any other iterative procedure.

The relaxation factor  $\omega$  in equation (11) is given in the form of the Chebyshev polynomials, but for actual computation it is more convenient to express  $\omega$  as

$$\left. \begin{aligned} \omega_{m+1} &= \frac{1}{1 - \frac{1}{4} \left[ \bar{\mu}^2(D) \omega_m \right]} && \text{for } m \geq 2 \\ \omega_1 &= 1 \\ \omega_2 &= \frac{2}{2 - \bar{\mu}^2(D)} \end{aligned} \right\} \quad (12)$$

#### Spectral Radius of Matrix D

As was pointed out in reference 1, it is very important to choose the relaxation factor  $\omega$  of equation (11) with great care in order to obtain an optimum rate of convergence of that equation. It is evident from equation (12) that  $\omega$  is a function of the spectral radius  $\bar{\mu}(D)$ .

If  $\underline{u}$  is any vector with positive components  $u_i$ , determination of non-trivial upper and lower bound estimates for the spectral radius of  $D$  may be obtained by applying the minimax theorem (ref. 4), which for the  $i^{\text{th}}$  iteration is given by

$$\min_i \left( \frac{\sum_j d_{i,j} u_j}{u_i} \right) \leq \bar{\mu}(D) \leq \max_i \left( \frac{\sum_j d_{i,j} u_j}{u_i} \right)$$

Since  $D$  is irreducible,  $\bar{\mu}(D)$  can be expressed finally as

$$\max_{\underline{u} \in R} \left[ \min_i \left( \frac{\sum_j d_{i,j} u_j}{u_i} \right) \right] = \bar{\mu}(D) = \min_{\underline{u} \in R} \left[ \max_i \left( \frac{\sum_j d_{i,j} u_j}{u_i} \right) \right] \quad (13)$$

#### The Iterative Procedure

The numerical solution of the axially symmetric Poisson equation (eq. (3)) now follows a procedure similar to that outlined in reference 1. The principal difference is in the iteration-matrix  $D$  of equation (9), which was derived in the previous section. Equation (11) is solved first with no space charge to obtain the Laplacian potential distribution inside region  $R$ . From this potential

distribution and the equations of motion, ion trajectories are calculated and the first-order space-charge-density-distribution function  $f(w, r, \theta, z)$  of equation (3) is determined. There is an additional feature in this computer program for pointwise calculation of ion trajectories that is not included in reference 1. In the event a trajectory goes through an almost  $90^\circ$  bend, a switch in the program changes the direction of sweep (e.g., from  $r$  to  $z$ ) of the increments from station  $n$  to station  $n + 1$ . The accuracy of computation is improved by this method.

A series of successive approximations to the space-charge-density function  $f(w, r, \theta, z)$  together with the necessary precautionary checks, as discussed in reference 1, then gives the solution of the axially symmetric Poisson equation.

### NUMERICAL EXAMPLE

The example chosen to demonstrate the numerical method of solution of the axially symmetric Poisson equation is a hollow-cylindrical-beam ion thruster. The mathematical model that has been analyzed is shown in figure 1. The coordi-

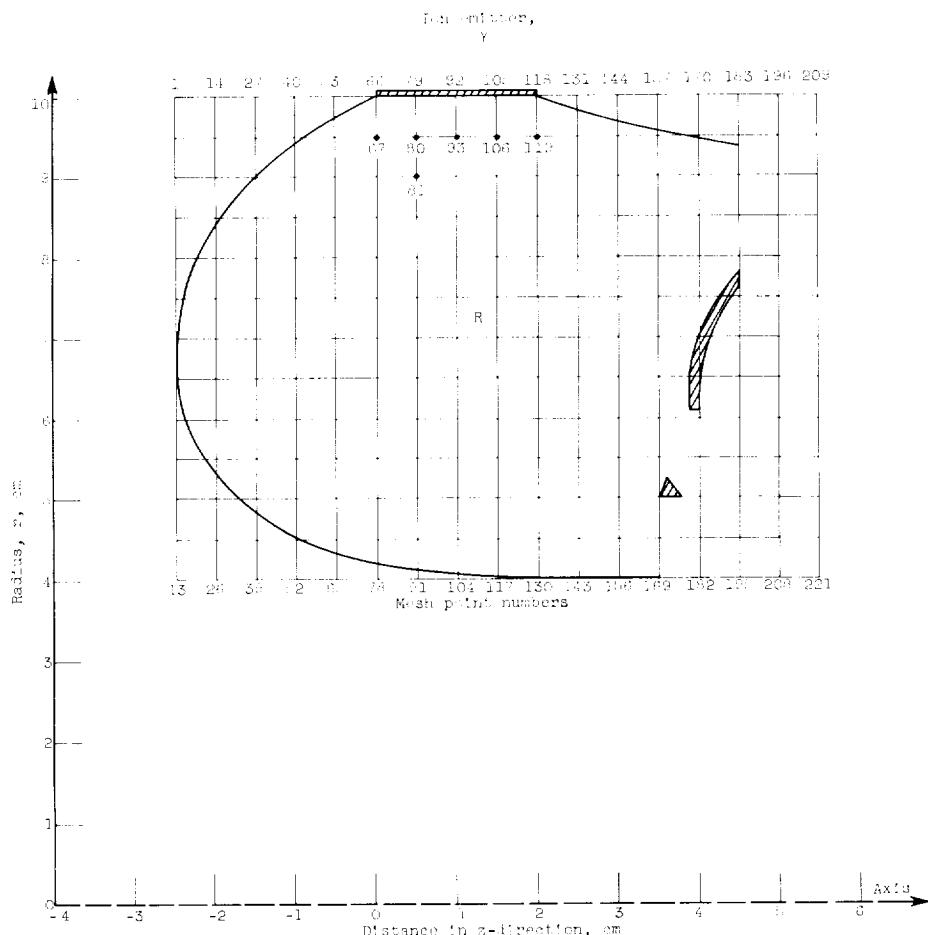
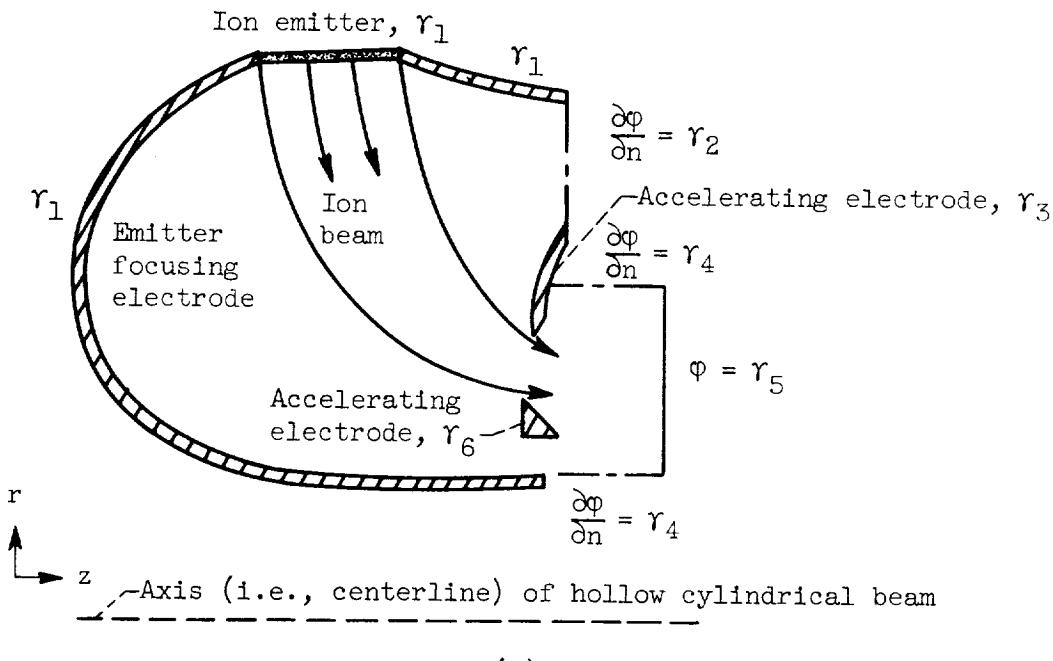


Figure 1. - Mathematical model for example problem from reference 3 for hollow-cylindrical ion beam.

nates of the electrode shapes were obtained from reference 3. The configuration was slightly modified to conform to physical reality, namely, the inner electrodes were assigned a definite thickness rather than lines as given in reference 3. The downstream external boundary conditions were chosen to permit the program to analyze in detail the aperture effect and the impingement current on the electrodes. These effects are of vital interest to good ion-thruster design. An IBM 7090 computer was used to solve the numerical example. (The computer program is given in appendix C.)

Operation of the hollow-cylindrical beam ion thruster is as follows: Ions are formed on the ion emitter, for example, the porous tungsten type, which is heated and at a positive potential relative to ground. The accelerator or inner electrodes are at a lower potential than the ion emitter, usually at a negative potential to prevent electrons from entering the thruster. With the assumption that an adequate flow of propellant is supplied (cesium vapor in this case), the potential field created between the ion emitter and the accelerator electrodes gives rise to space-charge-limited flow of the ions.

The boundaries were chosen as follows (see sketch (c)): (1) the ion emitter and the emitter focusing electrode  $r_1$  at a uniform potential of 2828.6 volts, (2) the portion of the external boundary between the emitter focusing electrode and the accelerator electrode  $r_2$  with the normal derivative equal to zero, (3) the accelerating electrode  $r_3$  at -1592.4 volts, (4) the portion of external



(c)

boundaries behind the accelerating electrodes  $r_4$  arbitrarily chosen with the normal derivative equal to zero, (5) the problem arbitrarily terminated with

ground potential  $\gamma_5$ , (6) the inner accelerating electrode  $\gamma_6$  at ground potential. All potentials are referenced to ground. The potentials  $\gamma_1$ ,  $\gamma_3$ , and  $\gamma_6$  resulted from inverting the potentials for electron flow given in reference 3. In this inversion,  $\gamma_6$  was arbitrarily chosen at ground potential.

The treatment and the location of the downstream boundaries  $\gamma_4$  and  $\gamma_5$  are arbitrary from the mathematical viewpoint. For a real ion thruster, however, these boundary conditions may be of considerable importance. As discussed in reference 1, various possibilities may arise that are primarily related to the problem of beam neutralization. It is not the purpose of this report to show the dependence of the selection of the downstream boundaries on the overall solution of the problem, but the reader is cautioned to use his own judgment or available experimental data to determine the approximate downstream boundary conditions for a given configuration. It was felt that the selection of the downstream boundary as shown in sketch (c) may be a good approximation to the actual operational conditions of this thruster.

Equipotentials of the solution of the Poisson equation for region R (shown in fig. 1) together with ion trajectories are shown in figure 2. The average current density of the ion beam at the emitter was calculated to be 0.147 ampere per square meter for the potential distribution shown in figure 2. This value is very low because of the "accel length" of 4 to 5 centimeters instead of 4 to 5 millimeters as may occur in real ion-thruster practice. From examination of the ion trajectories, the impingement current on the accelerator electrodes was calculated to be 5.0 percent of the ion emitter current.

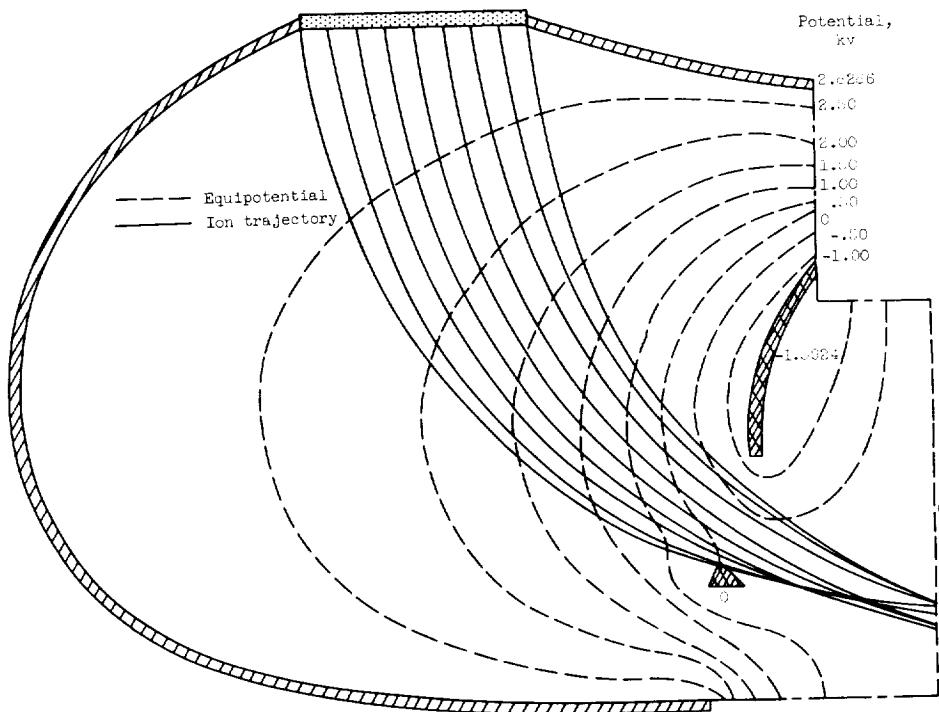


Figure 2. - Equipotentials and ion trajectories for space-charge-limited flow in region R of figure 1.

It is of interest to analyze the right-hand side (RHS) of equation (3), as was done for the two-dimensional case in reference 1, for the first row (or column) from the ion emitter (i.e., mesh points 67, 80, 93, 106, 119 in the first row from the emitter in fig. 1). If it is assumed that the trajectories very close to the emitter may follow straight lines, the RHS for the first row from the ion emitter can be written as

$$f \approx \frac{1}{\epsilon_0} \frac{j_E}{\bar{v}} \quad (14)$$

where  $j_E$  is obtained from the Child-Langmuir formula:

$$j_E = \frac{4}{9} \epsilon_0 \sqrt{\frac{2q}{m}} \frac{\Delta w^{3/2}}{h^2}$$

From the conservation of energy equation,

$$\bar{v} = \bar{v}_r = \sqrt{\frac{2q}{m} \Delta r_w}$$

If the potential along the first row is assumed almost constant (trajectories are straight lines),

$$\Delta w = \Delta r_w$$

and equation (14) becomes

$$f = \frac{4}{9} \frac{\Delta r_w}{h^2} \quad (15)$$

Substitution of equation (15) into equation (4) written, for example, for mesh point 80 (see fig. 1) gives

$$r(4w_{80} - w_{93} - w_{67} - w_{81} - r_1) + \frac{h}{2}(w_{81} - r_1) = \frac{4r}{9} \Delta r_w$$

Thus, when the potential distribution from the solution of the Laplacian equation is known, a quick check of the RHS as a first approximation of the Poisson equation at the first row can be made by multiplying the potential difference between the emitter and the mesh point of the first row (or column) by  $(4/9)r$ .

It should be noted that the electrode shapes obtained from reference 3 were calculated by a numerical method for a hollow beam described by an analytical solution in which the voltage and electric fields vary exponentially with respect to distance along the axis. The analytic solution of a hollow beam with cylindrical emitter did not take into account the aperture effect. Therefore, it would be expected that the electrode shapes calculated from that analytical solution may be in error for an actual ion-thrustor design.

The calculated impingement current of 5 percent herein might be attributed to accounting for the aperture effect as well as a change of the electric field around the electrodes resulting from introduction of a finite thickness to the electrodes. It is felt, therefore, that the agreement of this numerical method with the method of reference 3 for determination of electrode shapes for axially symmetric ion thrusters is good. Improvement in the ion optics may possibly be obtained by blocking segments of the ion emitter that had contributed the portion of the space-charge flow that was intercepted on the accelerator electrode. This type of improvement of ion optics is discussed in reference 2.

#### CONCLUDING REMARKS

The objective of this paper was to develop a numerical method of solution of the axially symmetric Poisson equation for mixed boundary conditions. The region for which the solution was sought was overlayed with mesh points, and for each mesh point the differential equation was replaced by finite-difference approximations. These approximations gave rise to a set of linear algebraic equations. For  $N$  mesh points in the region  $R$ , there were  $N$  linear algebraic equations with  $N$  unknowns, which resulted in an  $N$  by  $N$  real symmetric matrix - a Stieltjes matrix.

The theory of regular splittings was used because of the nonuniform main diagonal entries in the matrix. The matrix equation was solved by the Cyclic Chebyshev Semi-Iterative method on an IBM 7090 computer. To obtain fast convergence in the iterative process, an estimate of an optimum relaxation factor was used.

The method of solution presented herein is general for any type of external boundary that satisfies equation (2). The application of this method of solution to the analysis of the space-charge flow in a particular axially symmetric ion-thruster configuration was demonstrated. Agreement between ion trajectories obtained by this method and those obtained from an analytical solution was very good. The present method is superior to previous analytical methods in that the present method includes full account of the exhaust aperture that exists in real ion thrusters.

The method of solution presented herein may find use as a tool for diagnostic purposes by those working in this area. This method together with the two-dimensional method of analysis previously reported in references 1 and 2 extends the numerical program to include most of the ion-accelerator geometries currently being investigated.

Lewis Research Center  
National Aeronautics and Space Administration  
Cleveland, Ohio, February 4, 1963

## APPENDIX A

### SYMBOLS

$A$	matrix of matrix equation (eq. (6))
$A^{-1}$	matrix, inverse of matrix $A$
$B$	block diagonal matrix (eq. (7))
$B^{-1}$	matrix, inverse of matrix $B$
$b_{i,i}$	entries of matrix $B$
$C$	off-diagonal block matrix (eq. (7))
$c_{i,j}$	entries of matrix $C$
$D$	iteration matrix
$D_{1,2}$	matrix associated with odd- and even-number lines of $D$ , respectively
$d_{i,j}$	entries of matrix $D$
$f$	space-charge-density-distribution function for discrete case, v/sq m
$\underline{g}$	column vector, $B^{-1}\underline{k}$
$\underline{g}_{1,2}$	column vectors for odd- and even-number lines, respectively
$h$	mesh spacing, m
$j$	current density, amp/sq m
$\underline{k}$	column vector of matrix equation
$m$	particle mass, kg
$N$	number of mesh points in region $R$
$n$	outward normal
$q$	unit charge, coulombs
$R$	region in sketch (a)
$r$	subregion of $R$ ; cylindrical coordinate
$u$	component of $\underline{u}$

<u>u</u>	arbitrary column vector
<u>v</u>	average velocity, m/sec
w	potential-distribution function for discrete case, v
<u>w</u>	column vector of matrix equation
<u>w</u> <sub>1,2</sub>	column vectors for odd and even lines, respectively
z	cylindrical coordinate
$\alpha, \beta$	integers, (1 or 0)
$\Gamma$	external boundary of R
$\gamma$	discrete portion of external boundary
$\Delta$	increment or difference
$\delta$	increment of length in sketch (g) and eq. (B10)
$\nabla^2$	Laplacian operator
$\epsilon$	mathematical symbol representing "belongs to the set"
$\epsilon_0$	permittivity of free space, coulombs/(v)(m)
$\theta$	cylindrical coordinate
$\bar{\mu}(D)$	spectral radius of matrix D
$\rho$	space-charge-density-distribution function for continuous case, coulombs/cu m
$\sigma$	surface area bounded by subvolume $\tau$ in sketches (b) and (d)
$\tau$	subvolume of bounded space shown in sketches (b) and (d)
$\phi$	potential-distribution function for continuous case, v
$\omega$	relaxation factor

Subscripts:

E	emitter
i,j	number, 1, 2, . . . , N
m	number of iteration

N number of mesh points

r,z direction

Superscripts:

m number of iteration

o initial guess

l first iteration

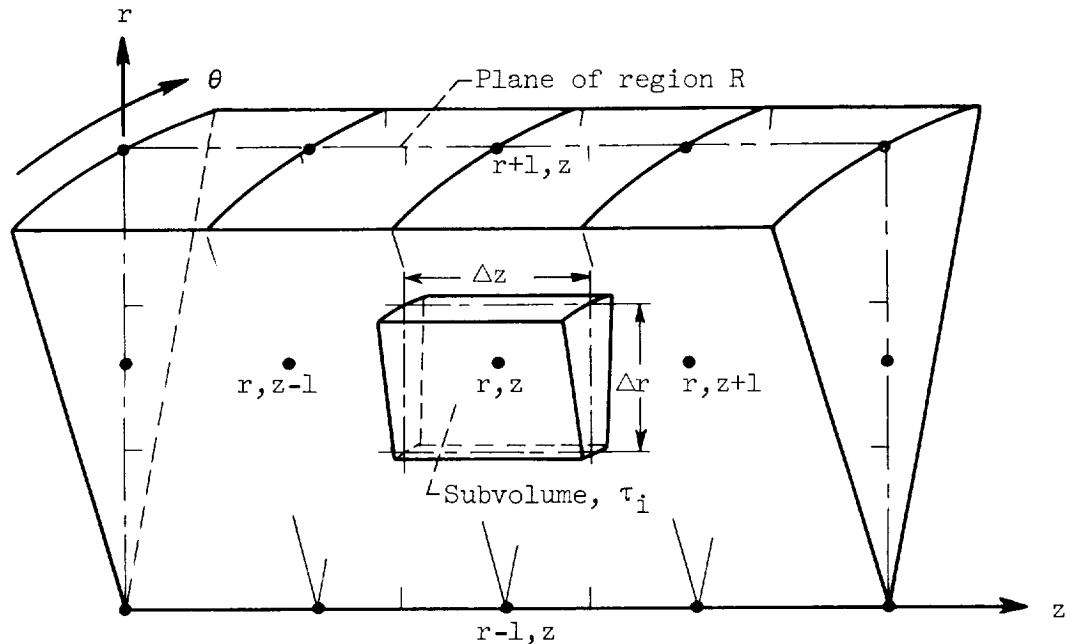
## APPENDIX B

### DERIVATION OF FINITE-DIFFERENCE EQUATIONS

#### ACCORDING TO REFERENCE 4

The Poisson equation in cylindrical coordinates is

$$-\nabla^2 w(r, \theta, z) = f(w, r, \theta, z) \quad (B1)$$



(d)

From the integration of equation (B1) over the subvolume  $\tau_i$  shown in sketch (d) it follows that

$$-\iiint_{\tau_i} \nabla^2 w(r, \theta, z) r \, dr \, d\theta \, dz = \iiint_{\tau_i} f(w, r, \theta, z) r \, dr \, d\theta \, dz \quad (B2)$$

By Green's theorem, the term on the left-hand side of equation (B2) can be reduced to a surface integral over the area  $\sigma_i$  bounding the subvolume  $\tau_i$ , and equation (B2) can be written as

$$-\iint_{\sigma_i} \frac{\partial w(r, \theta, z)}{\partial n} \, d\sigma = \iiint_{\tau_i} f(w, r, \theta, z) r \, dr \, d\theta \, dz \quad (B3)$$

where  $\partial w(r, \theta, z)/\partial n$  is the derivative in the direction of the outward normal to  $\sigma_i$ .

To obtain a five-point-formula approximation, the following numerical approximations to the integrals of equation (B3) are made. The function  $f(w, r, \theta, z)$  is assumed to be constant for the subvolume  $\tau_i$ , and therefore the RHS of equation (B3) becomes

$$\iiint_{\tau_i} f(w, r, \theta, z) r dr d\theta dz \approx f_i \iiint_{\tau_i} r dr d\theta dz \quad (B4)$$

For the axially symmetric case,  $w$  is independent of  $\theta$  and the normal derivatives of the left-hand side of equation (B3) are approximated by the central difference formula as

$$\frac{\partial w}{\partial z} \left( r, z + \frac{1}{2} \right) \approx \frac{w(r, z+1) - w(r, z)}{\Delta z} \quad (B5)$$

for the  $z$ -direction shown in sketch (d). Substitution of the preceding approximations into equation (B3) and integration gives the five-point formula, which can be written as

$$\begin{aligned} & - \left\{ \frac{w_{r,z+1} - w_{r,z}}{\Delta z} \left[ \pi \left( r + \frac{\Delta r}{2} \right)^2 - \pi \left( r - \frac{\Delta r}{2} \right)^2 \right] \frac{\Delta \theta}{2\pi} + \frac{w_{r,z-1} - w_{r,z}}{\Delta z} \right. \\ & \quad \times \left[ \pi \left( r + \frac{\Delta r}{2} \right)^2 - \pi \left( r - \frac{\Delta r}{2} \right)^2 \right] \frac{\Delta \theta}{2\pi} + \frac{w_{r+1,z} - w_{r,z}}{\Delta r} \left[ \left( r + \frac{\Delta r}{2} \right) \Delta \theta \Delta z \right] \\ & \quad \left. + \frac{w_{r-1,z} - w_{r,z}}{\Delta r} \left( r - \frac{\Delta r}{2} \right) \Delta \theta \Delta z \right\} = f_{r,z} r \Delta r \Delta \theta \Delta z \end{aligned} \quad (B6)$$

To simplify the matrix coefficients, it is necessary to choose  $\Delta r = \Delta z = h$ . Equation (B6) then reduces to

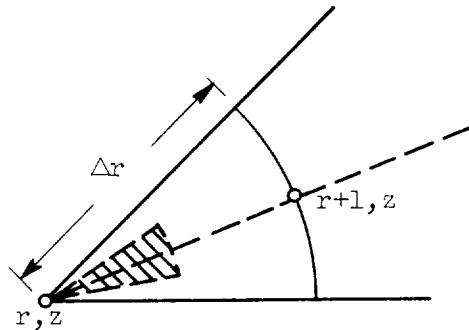
$$r \left( 4w_{r,z} - w_{r,z+1} - w_{r,z-1} - w_{r+1,z} - w_{r-1,z} \right) + \frac{h}{2} \left( w_{r-1,z} - w_{r+1,z} \right) = rh^2 f_{r,z} \quad (B7)$$

For  $r = 0$  (see sketch (e)), the five-point-formula approximation of equation (B3) is

$$\begin{aligned} & - \left\{ \frac{w_{r,z+1} - w_{r,z}}{\Delta z} \left[ \pi \left( \frac{\Delta r}{2} \right)^2 \frac{\Delta \theta}{2\pi} \right] + \frac{w_{r,z-1} - w_{r,z}}{\Delta z} \left[ \pi \left( \frac{\Delta r}{2} \right)^2 \frac{\Delta \theta}{2\pi} \right] \right. \\ & \quad \left. + \frac{w_{r+1,z} - w_{r,z}}{\Delta r} \left( \frac{\Delta r}{2} \Delta \theta \Delta z \right) \right\} = f_{r,z} \pi \left( \frac{\Delta r}{2} \right)^2 \frac{\Delta \theta}{2\pi} \Delta z \end{aligned}$$

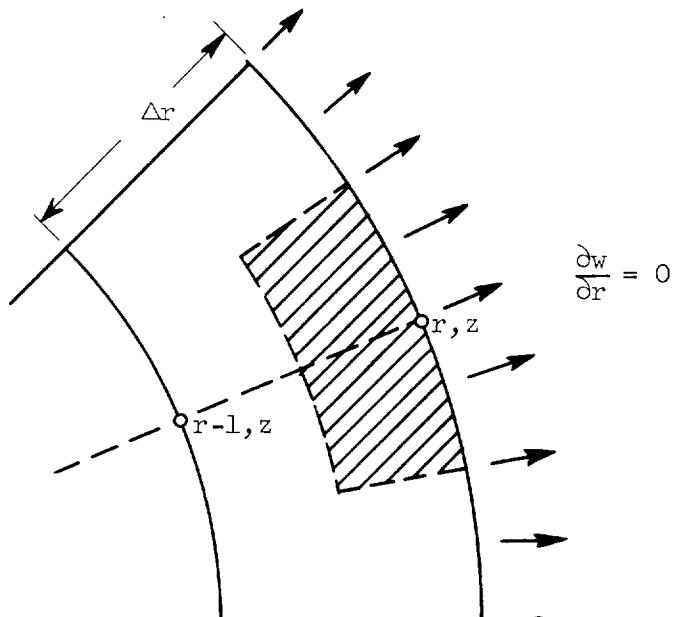
For  $\Delta r = \Delta z = h$ , it follows that

$$\frac{3w_{r,z} - w_{r,z+1} - w_{r,z-1}}{4} - \frac{w_{r+1,z} - w_{r-1,z}}{2} = \frac{h^2}{8} f_{r,z} \quad (B8)$$



(e)

For the case where the normal derivative is specified at the external boundary (see sketch (f)), the five-point-formula approximation of equation (B3) is



(f)

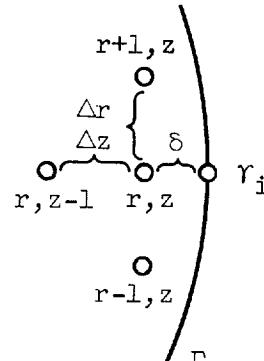
$$-\left\{ \frac{w_{r,z+1} - w_{r,z}}{\Delta z} \left[ \pi r^2 - \pi \left( r - \frac{\Delta r}{2} \right)^2 \right] \frac{\Delta \theta}{2\pi} + \frac{w_{r,z-1} - w_{r,z}}{\Delta z} \left[ \pi r^2 - \pi \left( r - \frac{\Delta r}{2} \right)^2 \right] \frac{\Delta \theta}{2\pi} \right. \\ \left. + \frac{w_{r-1,z} - w_{r,z}}{\Delta r} \left( r - \frac{\Delta r}{2} \right) \Delta \theta \Delta z \right\} = f_{r,z} \left[ \pi r^2 - \pi \left( r - \frac{\Delta r}{2} \right)^2 \right] \frac{\Delta \theta}{2\pi} \Delta z$$

For  $\Delta r = \Delta z = h$ , it follows that

$$\left(2r - \frac{3h}{4}\right)w_{r,z} - \left(\frac{r}{2} - \frac{h}{8}\right)w_{r,z+1} - \left(\frac{r}{2} - \frac{h}{8}\right)w_{r,z-1} - \left(r - \frac{h}{2}\right)w_{r-1,z} = f_{r,z} \frac{h^2}{2} \left(r - \frac{h}{4}\right)$$

(B9)

For regions with curved external boundaries (see sketch (g)) the numerical



(g)

approximation around the curved portion of the boundary  $\Gamma$  is as follows:

$$\begin{aligned} & - \left\{ \frac{w_{r+1,z} - w_{r,z}}{\Delta r} \left(r + \frac{\Delta r}{2}\right) \left(\frac{\Delta z + \delta}{2}\right) \Delta \theta + \frac{w_{r-1,z} - w_{r,z}}{\Delta r} \left(r - \frac{\Delta r}{2}\right) \left(\frac{\Delta z + \delta}{2}\right) \Delta \theta \right. \\ & \quad + \frac{r_i - w_{r,z}}{\delta} \left[ \pi \left(r + \frac{\Delta r}{2}\right)^2 - \pi \left(r - \frac{\Delta r}{2}\right)^2 \right] \frac{\Delta \theta}{2\pi} + \frac{w_{r,z-1} - w_{r,z}}{\Delta z} \\ & \quad \times \left. \left[ \pi \left(r + \frac{\Delta r}{2}\right)^2 - \pi \left(r - \frac{\Delta r}{2}\right)^2 \right] \frac{\Delta \theta}{2\pi} \right\} = f_{r,z} \left[ \pi \left(r + \frac{\Delta r}{2}\right)^2 - \pi \left(r - \frac{\Delta r}{2}\right)^2 \right] \frac{\Delta \theta}{2\pi} \left(\frac{\Delta z + \delta}{2}\right) \end{aligned}$$

For  $\Delta r = \Delta z = h$ , it reduces to

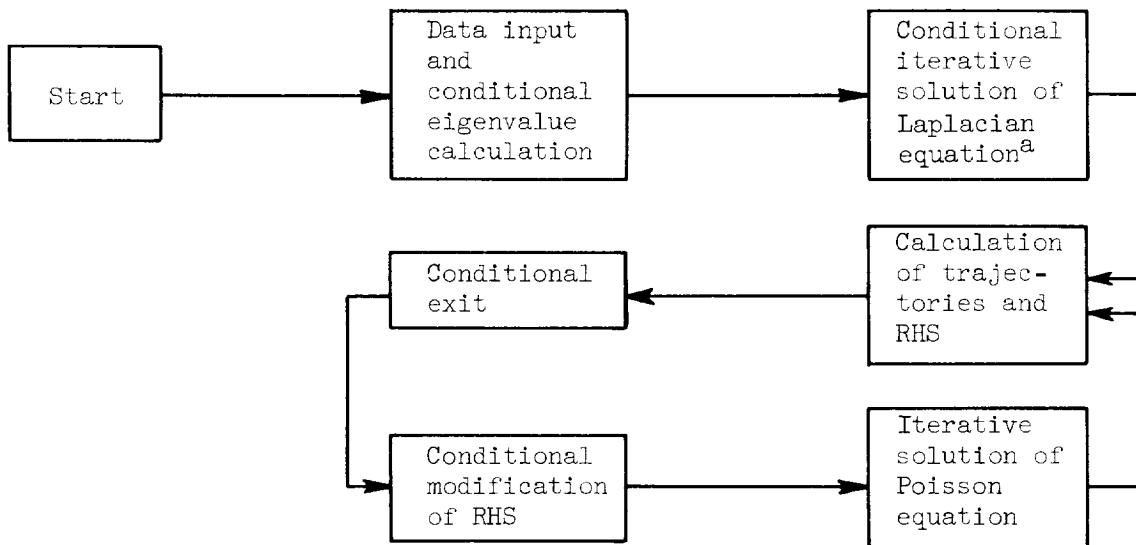
$$\begin{aligned} \frac{r}{\delta h} (h + \delta)^2 w_{r,z} - \left(r + \frac{h}{2}\right) \left(\frac{h + \delta}{2h}\right) w_{r+1,z} - \left(r - \frac{h}{2}\right) \left(\frac{h + \delta}{2h}\right) w_{r-1,z} \\ - \frac{rh}{\delta} r_i - rw_{r,z-1} = f_{r,z} rh \left(\frac{h + \delta}{2}\right) \end{aligned} \quad (\text{B10})$$

## APPENDIX C

### IBM 7090 ION-THRUSTOR FORTRAN CODE AND BLOCK DIAGRAM

By Carl D. Bogart

A schematic representation of the Fortran program is shown, followed by the symbol list, a flow chart of the program (fig. 3), and a complete Fortran listing with the data for a sample case.



<sup>a</sup>Same program as Poisson solution with RHS equal to zero.

### FORTRAN CODE SYMBOL LIST

#### Control Words:

ATX	x-emitter coordinates
ATY	y-emitter coordinates
IX0	positive, RHS prints out; negative or zero, RHS will not print out
JOT	number of lines of trajectories to be printed out
KAB	number of lines to traverse to obtain equipotential line
KAN	test point for RHS
KBA	test point for equipotential line if current is calculated from equal delta potential

KBN	vector to test for space-charge-limited flow
KCY	positive, all print-outs occur; negative or zero, no print-outs
KISW	positive, calculation of current is with equal delta-x; negative or zero, calculation of current is on equal delta potential
KRL	cycle counter
LINC	vector determining order of calculation of lines
MO	positive, test RHS upper bound; negative, test RHS lower bound; zero, no test
NCOOR	positive, emitter coordinates are beginning trajectory coordinates; negative or zero, beginning trajectory coordinates are calculated
NEM	number of emitter coordinates
NH	number of heading cards
NJOT	number of lines of trajectory output for Poisson solution
NLB	number of LB's to be read in
NLC	number of LC's to be read in
NLD	number of LD's to be read in
NLIN	number of lines in matrix equation
NPIT	negative, eigenvalue to be calculated; zero, potential input is from resistance paper; positive, potential input is from dump for restart
NPOL	first point in equipotential or first point for calculation of current from equal delta-x
NPONT	number of points per line
NRL	number of cycles
NSPAN	width of y-sweep if equipotential is calculated for horizontal emitter
NSWP	positive or negative, sweep is in y-direction; zero, sweep is in x-direction initially
NTJ	number of trajectories
NTOP	total number of points
NTP	number of KT's or XT's plus eight

NUL number of iterations on matrix equation  
NURL initially to change NUL for Poisson; later as a switch to indicate end of problem

Problem Specifications:

A atomic number of ions  
EPS convergence test for matrix equation  
H mesh size  
HGH lower coordinate of upper electrode  
HSL maximum slope test to change sweep  
JT vector of type numbers  
KT vector of relative subscripts  
LB vector defining first and last point to be calculated for each line  
LC vector determining trajectory and RHS calculation  
LD vector determining equipotentials calculation  
NJT number of JT's  
NKAN number of test points for over space-charge-limited flow plus one  
RX suppression factor for RHS  
SIZE step size in equipotential calculation  
VA emitter potential  
VAT upper potential in equipotentials calculation  
VB upper electrode potential  
VBT lower potential in equipotentials calculation  
VC lower electrode potential  
XLOW upper coordinate of lower electrode  
XLSL minimum slope test to change sweep  
XM error function term

XN	error function term
XQM	charge-to-mass ratio
XR	spectral radius of matrix
XT	vector of relative weights
YEP	permittivity of free space

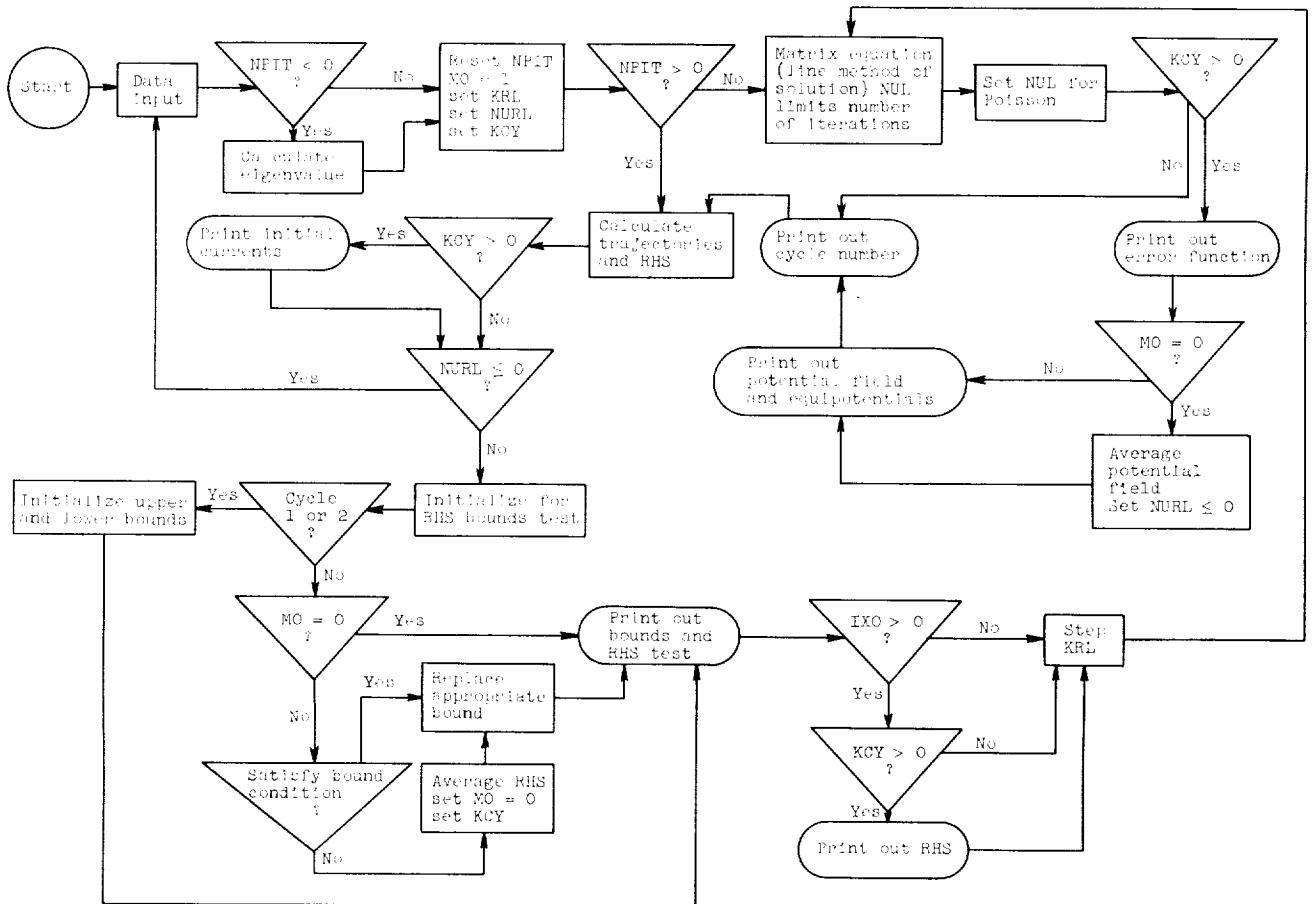


Figure 5. - Flow chart.

## ION-THRUSTOR FORTRAN CODE

```

      MAIN1
C MAIN ONE - DATA INPLT
      COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREG,NTP,
      1  NPII,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,LX,DELY,YEP,XCM,H,KISW,ETX,
      2  ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EPS,NPOL,
      3  NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RFUP,
      4  RFDOWN,XCL,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
      5  IXO,FCH,XLOW,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
      1  RF(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40)
      2  ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
      3  ETX(4C),ETY(4C),XCU(4C),KCY(1C)

200  MU=1
      READ INPUT TAPE 7,1CC,NH
      DO 13 J=1,NH
      READ INPUT TAPE 7,1C7
13   WRITE OUTPUT TAPE 6,107
      READ INPUT TAPE 7,1CC,NTOP,NPCNT,KAN,NSWF,NURL,NJCT,IXC
      READ INPUT TAPE 7,1CC,NLL,NLIN,NLB,NTP,NJT,NPIT
      READ INPUT TAPE 7,1CC,(LINC(J),J=1,NLIN)
      READ INPUT TAPE 7,1CC,(LB(J),J=1,NLB)
      CU 1 J=5,NTP,6
      K=J+5
      READ INPUT TAPE 7,1C1,(KT(M),M=J,K)
      1  READ INPUT TAPE 7,1C2,(XT(M),M=J,K)
      READ INPUT TAPE 7,1C3,(JT(M),M=1,NJT)
      READ INPUT TAPE 7,1C2,XR
      IF(NPIT) 2,3,23

C MATIN CALCULATES THE EIGENVALUE USED IN THE ITERATION SCHEME
2   CALL MATIN
3   JB=1
      DO 5 JL=1,NLIN
      READ INPUT TAPE 7,1CC,NC,NL,NR,NS
      READ INPUT TAPE 7,1C4,(LB(J),J=1,NC)
      NL=NL+LB(JB)
      NR=NR+LB(JB)
      CU 7 K=1,NC
      DO 4 JB=NL,NR
      4   U(JB)=UB(K)
      NL=NL+NS
      7   NR=NR+NS
      5   JB=JB+4
      CU TO 24
23   NTOP=NTOP
      CALL BCREAD(U(NTOP),L(1))
24   DO 6 J=1,NTOP
      UB(J)=U(J)
      6   RF(J)=0.
      READ INPUT TAPE 7,1CC,NCCCR,KBA,NTJ,NEM,NPCL,NPUL,KAB,NSPAN,JCT
      1  ,NLC,KISW,NRL,NKAN,NLD
      READ INPUT TAPE 7,1CC,(KBN(J),J=1,NKAN)
      READ INPUT TAPE 7,1CC,(LC(J),J=1,NLC)
      READ INPUT TAPE 7,1CC,(LD(J),J=1,NLD)
      READ INPUT TAPE 7,1C2,VA,VB,VC,H,EPS,RX
      READ INPUT TAPE 7,1C2,HSL,XLSL,HGH,XLCW,XN,XM

```

```

XMPR=2.*(XN*XN)**2/(XN**2+XM**2)
READ INPUT TAPE 7,1C5,YEP,XQM,A
READ INPUT TAPE 7,1C6,VAT,VBT,SIZE
XQM=XQM/A
READ INPUT TAPE 7,1C4,(ATX(J),J=1,NEM)
READ INPUT TAPE 7,1C4,(ATY(J),J=1,NEM)
IF(NCOOR) 1C,10,11
11 DO 12 J=1,NEM
   ETX(J)=ATX(J)
12 ETY(J)=ATY(J)
10 CONTINUE
KRL=NRL
READ INPUT TAPE 7,1CC,NPIT
READ INPUT TAPE 7,1CC,(KCY(J),J=1,14)
IF(NPIT) 25,25,26
C UCAL SOLVES THE MATRIX EQUATION
25 CALL UCAL
C MNTRI CALCULATES THE RHS AND TRAJECTORIES
26 CALL MNTRI
GO TO 2C0
100 FORMAT(14I5)
101 FORMAT(6I5)
102 FORMAT(6F10.5)
103 FORMAT(13I5)
104 FORMAT(7F10.5)
105 FORMAT(7E10.5)
106 FORMAT(3F10.5)
107 FORMAT(72H
1
END

```

```

SUBROUTINE MATIN
C MATIN CALCULATES THE EIGENVALUE
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,CX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NIJ,KBA,VA,VB,VC,ACCCR,KCH,HSL,XLSL,EPS,NFCL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LC,VAT,VBT,SIZE,KBN,RFUP,
4 RFDOWN,XCL,NOT,MC,KCY,NSKP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5 IXO,HGH,XLOW,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(1C),LRH(4000),UB(4000),JT(4000),
1 RH(4CCC),L(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCU(4C),KCY(10)
      DIMENSION A(63),X(21)
      EQUIVALENCE (LRH(1),A(1)),(LRH(100),X(1))

C MIT=NUMBER OF ITERATIONS ON EIGENVALUE CALCULATION
C KWR=- OR 0,NU INTERMEDIATE OUTPUT
C KWR=+,PRINT INTERMEDIATE ITERATIONS ON EIGENVALUE CALCULATION
      READ INPUT TAPE 7,102,MIT,KWR
C COLUMN VECTOR INITIALIZED
50      DO 46 J=1,NTOP
        IF(JT(J)>43,43,44
43      U(J)=C.
        GO TO 46
44      U(J)=1.
46      CONTINUE
      JS=-1
C ITERATIVE LOOP
      DO 41 M=1,MIT
        JB=1
C LOOP ON NUMBER OF LINES
      DO 31 ND=1,NLIN
        IF(LINC(ND)) 23,31,23
23      KA=LB(JB)
        KB=LB(JB+1)+KA
        KC=LB(JB+2)+KA
        JV=1
        JL=1
        DO 13 K=KB,KC
          SLM=C.
          JZ=JT(K)
          IF(JZ) 13,13,14
14      DO 15 JY=3,5
          JX=JZ+JY
          A(JL)=XT(JX)
15      JL=JL+1
        DO 16 JY=1,2
          JX=JZ+JY
          Jw=K+KT(JX)
16      SLM=SLM+XT(JX)*L(Jw)
          X(JV)=SLM
          JV=JV+1
13      CONTINUE
        N=KC-KB+1
        CALL MATRIX(N,A(2),X)
C MATRIX IS USED FOR THE FORWARD AND BACK SUBSTITUTION
        JV=1

```

```

      CO 17 J=KB,KC
      RH(J)=X(JV)
17    JV=JV+1
31    JB=JB+4
C SWITCH ALLOWING MATRIX TO BE APPLIED TWICE
      IF(JS) 32,34,34
32    DO 33 J=1,NTOP
      LB(J)=U(J)
33    U(J)=RH(J)
      GO TO 41
C DETERMINATION OF SMALLEST AND LARGEST RATIOS FOR EIGENVALUE
34    XL=C.
      XS=1.
      DO 39 JD=1,NTOP
      IF(U(JD))39,39,35
35    X=RH(JD)/LB(JD)
      IF(XL-X)36,37,37
36    XL=X
      NNL=JD
37    IF(XS-X)39,39,38
38    XS=X
      NS=JD
39    CONTINUE
      IF(KWR) 51,51,52
52    WRITE OUTPUT TAPE 6,1C1,M,XS,XL,NS,NNL
51    YL=RH(NNL)
      DO 40 JD=1,NTOP
40    U(JD)=RF(JD)/YL
      IF(XL-XS-1.0E-07)42,42,41
41    JS=-JS
42    XR=SQRT(.5*(XL+XS))
      WRITE OUTPUT TAPE 6,1C3,XR
47    RETURN
101   FORMAT(2OH     LOW     HIGH        I5,2F13.8,5H      216)
102   FORMAT(14I5)
103   FORMAT(4HCXR= F1C.8)
END

```

```

SUBROUTINE MATRIX(N,A,X)
C MATRIX IS USED FOR THE FORWARD AND BACK SUBSTITUTION
COMMON LRH,UB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREG,NTP,
1  NPIT,XEP,NXEP,CL,APCNT,AX,AY,VX,VY,DY,DELY,YEF,XCM,F,KISW,ETX,
2  ETY,PTY,PTX,NAJ,NIJ,KBA,VA,VB,VC,NCLCR,KCH,HSL,XLSL,EFS,NFCL,
3  NPL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUF,
4  RFDOWN,XCU,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5  IXU,FGH,XLOW,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(10),LRH(4000),UB(4000),JT(4000),
1  RF(4CCC),L(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2  ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3  ETX(4C),ETY(4C),XCL(4C),KCY(10)
      DIMENSION A(63),X(21)

M=3*N-2
A(M+1)=C.
A(2)=A(2)/A
X=X/A
      IF(N-1) 12,12,9
9   K=2
DO 1C J=4,M,3
A(J)=A(J)-A(J-1)*A(J-2)
A(J+1)=A(J+1)/A(J)
X(K)=(X(K)-A(J-1)*X(K-1))/A(J)
10  K=K+1
      K=K-1
DO 11 J=1,M,3
NB=M-J-1
K=K-1
11  X(K)=X(K)-A(NB)*X(K+1)
12  RETURN
END

```

```

SUBROUTINE UCAL
C UCAL SOLVES THE MATRIX EQUATION
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,CX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,P TY,P TX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EPS,NFCL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBA,RHUP,
4 RFDOWN,XCU,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5 IXO,FGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),LRH(4000),UB(4000),JT(4000),
1 RF(4CCC),U(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCL(4C),KCY(10)
DIMENSION A(63),X(21)
EQUIVALENCE(LRH(1),A(1)),(LRH(100),X(1))
SX=1.
46 XW=1.
CO 30 NLU=1,NLL
NEZ=NLU
XEP=0.
XM=.25*XR**2
NOD=1
38 JB=1
CO 29 NL=1,NLIN
IF(LINC(NL)) 35,35,33
33 KA=LB(JB)
KB=LB(JB+1) +KA
KC=LB(JB+2) +KA
JV=1
JL=1
CO 28 K=KB,KC
SUM=0
JZ=JT(K)
IF(JZ) 28,28,27
27 CO 22 JY=3,5
JX=JZ+JY
A(JL)=XT(JX)
22 JL=JL+1
CO 26 JY=1,5
JX=JZ+JY
IF(KT(JX)) 34,26,34
34 IF(XT(JX)) 26,26,25
25 JW=K+KT(JX)
SUM=SLM+XT(JX)*L(JW)
26 CONTINUE
X(JV)=SLM+XT(JZ)*RH(K)
JV=JV+1
28 CONTINUE
N=KC-KB+1
C MATRIX IS USED FOR THE FORWARD AND BACK SUBSTITUTION
CALL MATRIX (N,A(2),X)
JV=1
CO 24 K=KB,KC
IF(JT(K)) 24,24,23
23 CIF=X(JV)-U(K)
JV=JV+1
C MATRIX EQUATION

```

```

U(K)=XW*DIF+U(K)
DIF=ABSF(DIF)
IF(DIF-XEP) 24,24,21
21 XEP=DIF
NXEP=K
24 CONTINUE
JB=JB+4
29 LINC(NL)==LINC(NL)
IF(XW-1.) 40,31,40
40 XW=1./(1.-XM*XW)
GO TO 41
31 XW=1./(1.-2.*XM)
41 IF(NOD) 37,37,36
36 NOD=C
GO TO 38
37 XCON=XEP*XMPR
IF(XCON -EPS) 42,42,30
30 CONTINUE
42 KNUT=KBN(1)
J=2
DO 43 K=1,KNUT
JN=KBN(J)
IF(U(JN)-VA) 43,44,44
43 J=J+1
GO TO 47
44 SX=SX*RX
WRITE OUTPUT TAPE 6,1C0,JN,L(JN),SX
DO 45 J=1,NTOP
U(J)=UB(J)
45 RF(J)=RH(J)*RX
KAN=KAN
RFUP=RH(KAN)
GO TO 46
47 IF(NRL-KRL) 2,1,2
1 NUL=NRL
NTOP=NTOP
CALL BCDUMP (U(NTOP),L(1))
2 IWRL=NRL-KRL+XABSF(MC)
IF(KCY(IWRL)) 4,4,3
C TWOUT PRINTS OUT THE POTENTIAL FIELD
3 CALL TWOUT(NEZ)
C EQLINE CALCULATES THE EQUIPOTENTIALS
CALL EQLINE
4 WRITE OUTPUT TAPE 6,1C1,IWRL
RETURN
100 FORMAT(31HPOINT/VALLE/TOTAL SUPRESSION I5,2F15.5)
1C1 FORMAT(8HCR LOOP I2)
END

```

```

        SUBROUTINE EQLINE
C EQLINE SOLVES FOR THE EQUIPOTENTIALS
      COMMON LRH,L8,J1,RH,L,XT,KT,LINC,LB,XR,ATCP,NUL,ALIN,NREC,NTP,
     1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DLY,YEP,XGM,F,KISW,ETX,
     2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EFS,NPOL,
     3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBA,RFUP,
     4 RFDOWN,XCL,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
     5 IXU,FGH,XLUH,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(1C),LRH(4000),UB(4000),JT(4000),
     1 RF(4CCC),U(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
     2 ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(40),ATY(4C),
     3 ETX(4C),ETY(4C),XCL(4C),KCY(10)
      JB=1
      IF(SIZE) 1,27,1
1      POTEN=VAT
2      JE=LD(JB)
3      JD=LD(JB+1)-1
4      JC=LD(JB+2)
5      DX=LD(JB+3)
6      DX=DX*H
7      BX=C.
8      JED=JE+JD
9      L=1
10     IF(JL-1) 3,3,4
11     AX=C.
12     DO 22 JJ=1,JC
13     KS=1
14     AAY=C.
15     DO 19 K=JE,JED
16     IF(KS) 8,8,7
17     M=1
18     J=K-LD(JB+1)
19     GO TO 9
20     J=K-1
21     IF(JT(K)+JT(J)) 18,18,28
22     IF((U(K)-POTEN)*(L(J)-PCTEN)) 10,10,18
23     DIF=ABSF(L(J)-U(K))
24     IF(DIF) 13,13,11
25     IF(M) 12,14,12
26     VX(L)=ABSF(L(J)-POTEN)/DIF*DX+AX
27     VY(L)=AAY
28     GO TO 15
29     VX(L)=AX+DX
30     VY(L)=AAY
31     GO TO 15
32     VX(L)=AX+DX
33     VY(L)= ABSF(L(J)-POTEN)/DIF*DX+AAY
34     IF(L-7) 17,16,16
35     WRITE OUTPUT TAPE 6,100,PCTEN,(VX(I),VY(I),I=1,7)
36     L=C
37     L=L+1
38     AAY=AAY+DX
39     CONTINUE
40     IF(KS) 21,21,20
41     KS=C
42     M=C

```

```
JE=JE+1
GO TO 5
21 JE=JE+JD
JED=JE+JD
BX=BX+DX
AX=AX+DX
22 CONTINUE
IF(L-2) 25,23,23
23 DO 24 J=L,7
VX(J)=0.
24 VY(J)=0.
WRITE OUTPUT TAPE 6,1CO,PCTEN,(VX(I),VY(I),I=1,7)
25 POTEN=POTEN-SIZE
IF(POTEN-VBT) 27,26,26
26 AX=AX-BX
GO TO 2
27 RETURN
1CO FORMAT(17HCPOTENTIAL (X,Y) F8.1,2H 7(2H (F5.3,1F,F5.3,2F) ))
END
```

```

SUBROUTINE TWOUT(KK)
C TWOUT PRINTS OUT THE POTENTIAL FIELD
COMMON LRH,LB,J1,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREG,NTP,
1 NPIT,XEP,NXEP,CL,NPONT,AX,AY,VX,VY,DLY,YEP,XQM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EPS,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBA,RHUP,
4 RFDOWN,XCL,NUT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5 IXO,HGH,XLOW,XMPR,NEM
      DIMENSION KBN(20),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1 RF(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40)
2 ,VX(40),KCH(40),PTX(40),PTY(40),AY(40),AX(40),ATX(40),ATY(40),
3 ETX(40),ETY(40),XCU(40),KCY(10)
NN=KK
      WRITE OUTPUT TAPE 6,1C2,NN
      XCON=XEP*XMPR
      WRITE OUTPUT TAPE 6,1C1,NXEP,XCCN,EPS
      IF(MO) 10,3,1C
1C      K=1
      DO 1 J=1,NTOP
      KT(K)=J
      XT(K)=U(J)
      K=K+1
      IF(K-5) 1,2,2
2      WRITE OUTPUT TAPE 6,1C0,(KT(L),L=1,8),(XT(L),L=1,8)
      K=1
      CONTINUE
      IF(K-2) 8,6,6
6      CU 7 J=K,8
      KT(J)=C
      XT(J)=C.
      WRITE OUTPUT TAPE 6,1C0,(KT(L),L=1,8),(XT(L),L=1,8)
8      RETURN
3      WRITE OUTPUT TAPE 6,1C4
      DO 11 J=1,NTOP
11      L(J)=.5*(L(J)+LB(J))
      JOT=NJOT
      NURL=-7777
      GO TO 1C
100     FORMAT(1H 8I5,8F10.3)
101     FORMAT(21HPOINT/ERROR/EPISILCN    I5,2F15.6)
102     FORMAT(13HOL-ITERATION   I3)
104     FORMAT(23HOL-FIELD IS AVERAGED.      )
      END

```

```

SUBROUTINE MNTRI
C MNTRI COORDINATES THE RHS AND TRAJECTORY CALCULATION
COMMON URH,UB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
      NPIT,XEP,NXEP,CL,NPONT,AX,AY,VX,VY,DX,DELY,YEP,XCM,F,KISW,ETX,
      ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,ACCCR,KCH,HSL,XLSL,EPS,NPOL,
      NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VET,SIZE,KBN,RHUP,
      RFDOWN,XCU,NOT,MO,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
      IXO,HGH,XLOW,XMPR,NEM
      DIMENSION KBN(20),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
      RH(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40)
      ,VX(40),KCH(40),PTX(40),PTY(40),AY(40),AX(40),ATX(40),ATY(40),
      ETX(40),ETY(40),XCU(40),KCY(10)
      NPIT=C
301   DO 1 J=1,40
      PTX(J)=C.
      PTY(J)=C.
      VY(J)=.CCCCC1
      VX(J)=0.
      CU(J)=0.
1     KCH(J)=-1
      DO 35 J=1,NTOP
      LB(J)=U(J)
35     URH(J)=RH(J)
      NOT=JOT
      NAJ=NTJ-1
      IF(NCOOR) 2,2,3
C ARC CALCULATES THE Emitter COORDINATES IF NECESSARY
2     CALL ARC
      NCOOR=1
3     IF(KISW) 13,13,12
C PEQ CALCULATES THE EQUIPOTENTIAL LINE FOR CU CALCULATION
13    CALL PEQ
12    DO 11 J=1,NTJ
      AX(J)=ETX(J)
11     AY(J)=ETY(J)
      DX=LC(1)
      DX=DX*H
      DELY=DX
      JC=LC(2)
      JCX=LC(3)
      JE=LC(4)
      JEX=LC(5)
      JC=LC(6)
      JCX=LC(7)
      RM=DX/YEP
      AAX=C.
      NS=C
      JED=JD
      DO 22 JN=1,JC
      DO 7 J=1,NTJ
7       AY(J)=AY(J)+DX
      IF(NS) 4,4,5
C CURRENT CALCULATES THE CURRENT PER UNIT LENGTH
4     CALL CURRENT
      NS=1
C TRAJY CALCULATES THE TRAJECTORY COORDINATES IN THE Y SWEEP

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```

5      CALL TRAJY(JE,JED)
17      IF(NOT)6,6,17
17      NOT=NOT-1
17      WRITE OUTPUT TAPE 6,1C0,NCT,AY(1),(K,AX(K),VX(K),VY(K),K=1,NTJ)
6      DO 203 K=JE,JED,JDY
203      RF(K)=0.
C CALRY CALCULATES THE RHS IN THE Y SWEEP
      CALL CALRY(JE,JED)
      DO 21 K=JE,JED,JDY
21      RF(K)=RH(K)*RM
      BIG=0.
      SML=1.
      DO 34 J=1,NTJ
      S=ABSF(VY(J)/VX(J))
      IF(S-BIG) 33,33,36
36      BIG=S
33      IF(SML-S) 34,34,19
19      SML=S
34      CONTINUE
C SLOPE TEST FOR CHANGE OF SWEEP
      IF(BIG-FSL) 55,55,18
18      IF(SML-XSL) 55,55,23
23      JED=JED+1
22      JE=JE+1
      WRITE OUTPUT TAPE 6,1C4
      GO TO 307
55      JX=(AY+.CC01)/DX
      NOT=C
      DO 209 J=1,NAJ
209      CU(J)=-CU(J)
      JE=JEX+JX
      DO 200 JN=1,JCX
      JED=JEX+JDY-1
      AAX=AAX+DX
C TRAJX PICKS UP THE TRAJECTORIES IN THE X SWEEP
      CALL TRAJX(JEX,JDY,AAX)
C CORRCT CONDITIONALLY TERMINATES OR SHIFTS THE TRAJECTORIES
      CALL CORRCT(AAX,JN)
      IF(NOT) 8,8,9
9      NOT=NOT-1
      WRITE OUTPUT TAPE 6,1C1,NCT,AAX,(K,AY(K),VX(K),VY(K),K=1,NTJ)
8      JR=JE+1
      DO 202 K=JR,JED
202      RF(K)=0.
C CALRX CALCULATES THE RHS IN THE X SWEEP
      CALL CALRX(JEX,JED,JE)
      DO 201 K=JR,JED
201      RF(K)=RH(K)*RM
      JE=JE+JDY
200      JEX=JEX+JDY
307      JUT=C
      IWRL=NRL-KRL+1
      IF(KCY(IWRL)) 304,304,305
305      WRITE OUTPUT TAPE 6,102,(J,XCL(J),J=1,NAJ)
      SLM=0.
      DO 1C J=1,NAJ

```

```
1C      SLM=SLM+XCL(J)
       WRITE OUTPUT TAPE 6,1C3,SLM
304      IF(NURL) 303,3C3,3C2
C RTEST CHECKS THE UPPER AND LOWER BOUNDS ON RHS
302      CALL RTEST
C UCAL SOLVES THE MATRIX EQUATION
      CALL UCAL
      GO TO 301
303      RETURN
1C0      FORMAT(I5,F10.5,(1CH X VX VY   I5,3E15.6))
1C1      FORMAT(I5,F10.5,(1CH Y VX VY   I5,3E15.6))
1C2      FORMAT(17HOINITIAL CLRRENTS // (7(1H ,I2,E14.6)))
1C3      FORMAT(8HCTOTAL E14.6)
1C4      FORMAT(24HO TEST ON SLOPES UNMET.    )
      END
```

```

SUBROUTINE ARC
C ARC CALCULATES THE BEGINNING TRAJECTORY COORDINATES
COMMON LRH,LB,JI,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,AREG,NTP,
1 NPIT,XEP,NXEP,CL,APCNT,AX,AY,VX,VY,CX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XSL,EPS,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUP,
4 RHDOWN,XCU,NOT,MO,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5 IXO,HGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1 RH(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCU(4C),KCY(10)
SLM=C.
DO 1 J=2,NEM
1 SUM=SUM+SQR TF((ATX(J)-ATX(J-1))**2+(ATY(J)-ATY(J-1))**2)
ARCL=NTJ-1
ARCL=SUM/ARCL
K=1
REM=0.
J=1
HYP=ARCL
ETX(K)=ATX(K)
ETY(K)=ATY(K)
K=K+1
2 XDIFF=ATX(J)-ATX(J+1)
YDIF=ATY(J+1)-ATY(J)
SA=SQR TF(XDIF**2+YDIF**2)
CSA=YDIF/SA
SNA=XDIF/SA
3 ETX(K)=ATX(J)-HYP*SNA
ETY(K)=ATY(J)+HYP*CSA
HYP=ARCL+HYP
K=K+1
REM=REM+ARCL
IF(K-NTJ) 4,6,6
4 IF(SA-REM-ARCL) 5,3,3
5 J=J+1
HYP=ARCL-SA+REM
REM=HYP-ARCL
GO TO 2
6 CONTINUE
NEM=NEM
ETX(K)=ATX(NEM)
ETY(K)=ATY(NEM)
IF(JOT) 8,8,7
7 WRITE OUTPUT TAPE 6,1C3,SLM,ARCL
WRITE OUTPUT TAPE 6,1C0,(J,ATX(J),ATY(J),J=1,NEM)
WRITE OUTPUT TAPE 6,1C2,(J,ETX(J),ETY(J),J=1,NTJ)
8 CONTINUE
100 FORMAT (12HCX,Y-EMITTER//(7(1H,I2,2H (F5.3,1H,F5.3,2H) )))
102 FORMAT (16HCX,Y-BEGIN TRAJ./(7(1H,I2,2H (F5.3,1H,F5.3,2H) )))
103 FORMAT (30H ARC,DELTA ARC LENGTH 2F10.5)
RETURN
END

```

```

SUBROUTINE PEQ
C PEQ CALCULATES THE EQUIPOTENTIAL LINE FOR THE CURRENT DENSITY
C CALCULATION
      COMMON LRH,LB,JT,RH,L,XT,KI,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1     NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,CX,DELY,YEP,XCM,F,KISW,ETX,
2     ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EPS,NPOL,
3     NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VET,SIZE,KBN,RHUP,
4     RFDOWN,XCU,NOT,MO,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5     IXO,HGH,XLOW,XMPR,NEM
      DIMENSION KBN(20),LD(10),LC(10),LRH(4000),UB(4000),JT(4000),
1     RF(4000),U(4000),XT(630),KI(630),LINC(50),LB(200),CU(40),VY(40)
2     ,VX(40),KCH(40),PTX(40),PTY(40),AY(40),AX(40),ATX(40),ATY(40),
3     ETX(40),ETY(40),XCU(40),KCY(10)
      KBA=KBA
      POTEN=U(KBA)
      CX=H
      L=0
      KAB=KAB
      IF(NSWP) 32,24,32
32    AAY=C.
      JE=LC(4)
      JC=LC(6)
      DO 8 JJ=1,KAB
      AAX=C.
      JED=JE+NSPAN
      DO 7 K=JE,JED
      J=K-1
      IF((U(K)-POTEN)*(L(J)-POTEN)) 1,1,7
1     DIF=ABSF(U(J)-U(K))
      L=L+1
      PTX(L)=AAY
      IF(DIF) 2,6,2
2     PTY(L)=ABSF(U(J)-POTEN)/DIF*DX+AAX
      GO TO 7
6     PTY(L)=AAX+DX
7     AAX=AAX+DX
      AAY=AAY+DX
8     JE=JE+JD
      DO 11 J=1,L
      LL=L-J+1
      T=0.
      DO 10 I=1,LL
      IF(T-PTX(I)) 9,9,10
9     T=PTX(I)
      NN=I
10    CONTINUE
      PP=PTX(LL)
      PTX(LL)=T
      PTX(NN)=PP
      PP=PTY(LL)
      PTY(LL)=PTY(NN)
      PTY(NN)=PP
11    CONTINUE
      DO 14 J=2,L
      IF(PTX(J)-PTX(J-1)) 14,12,14
12    NN=L-1

```

```

DO 13 JJ=J,NN
PTX(JJ)=PTX(JJ+1)
PTY(JJ)=PTY(JJ+1)
13 CONTINUE
50 IF(JOT) 60,60,70
70 WRITIE OUTPUT TAPE 6,1CO,(J,PTX(J),PTY(J),J=1,L)
60 RETURN
24 AAX=C.
JE=LC(5)
JC=LC(6)
DO 28 JJ=1,KAB
AAY=0.
JED=JE+JD-1
DO 27 K=JE,JED
J=K-JD
IF((L(K)-POTEN)*(L(J)-POTEN)) 21,21,27
21 DIF=ABSF(L(J)-U(K))
L=L+1
PTY(L)=AAY
IF(DIF) 22,26,22
22 PTX(L)=ABSF(L(J)-POTEN)/DIF*DX+AAX
GO TO 27
26 PTX(L)=AAX+DX
27 AAY=AAY+DX
AAX=AAX+DX
28 JE=JE+JD
DO 41 J=1,L
LL=L-J+1
T=0.
JJ=LL
DO 40 I=1,LL
IF(T-PTY(I)) 39,39,40
39 T=PTY(I)
NN=I
40 CONTINUE
PP=PTY(JJ)
PTY(JJ)=T
PTY(NN)=PP
PP=PTX(JJ)
PTX(JJ)=PTX(NN)
PTX(NN)=PP
41 CONTINUE
DO 44 J=2,L
IF(PTY(J)-PTY(J-1)) 44,42,44
42 NN=L-1
DO 43 JJ=J,NN
PTX(JJ)=PTX(JJ+1)
43 PTY(JJ)=PTY(JJ+1)
44 CONTINUE
GO TO 50
100 FORMAT(18HCX,Y-EQUIPOTENTIAL//,(7(1H ,I2,2H (F5.3,1H,F5.3,2F) )))
END

```

```

      SUBROUTINE CURRNT
C CURRNT CALCULATES THE CURRENT PER UNIT LENGTH
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1  NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DX,DELY,YEP,XCM,F,KISW,ETX,
2  ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCR,KCH,HSL,XLSL,EPS,NPOL,
3  NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RFUP,
4  RFDOWN,XCL,NOT,NC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5  IX0,FGH,XLOw,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(1C),LRH(4000),UB(4000),JT(4000),
1  RF(4CC0),U(4CCC),XT(63C),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2  ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(4C),ATY(4C),
3  ETX(4C),ETY(4C),XCL(4C),KCY(1C)
      DIMENSION HX(2C)
      XK=2./9.*YEP*SQR TF(2.*XQM)*H
      IF(KISW) 20,2C,3C
20    DO 1C J=1,3C
      KBA=KBA
      DELL=VA-U(KBA)
      XK=XK*DELL*SQR TF(DELL)
      N=J
      K=(PTX(J)+.CCC1)/ETX
      IF(K) 1C,1C,11
10    CONTINUE
11    J=1
      L=1
12    XD=PTY(N)-ETY(J)
      YD=PTY(N+1)-ETY(J+2)
      XMID=.5*(XD+YD)
      UP=.5*(XMID+XD)
      DN=.5*(XMID+YD)
      FX(L)=UP
      FX(L+1)=DN
      L=L+2
      J=J+2
      N=N+1
      IF(L-NAJ) 12,13,13
13    DO 14 J=1,NAJ
14    CU(J)=      XK/HX(J)**2
40    DO 5C J=1,NAJ
50    XCU(J)=CL(J)
      RETURN
30    K=NPOL
      JC=NPUNT
      XK=XK/H**2
      J=1
39    JUB=K+JC
      LA=.5*(L(K)+L(JCB ))
      LL=.5*(LA+L(K))
      UA=.5*(LA+U(JUB ))
      DELL=VA-LL
      CL(J)=XK*DELL*SQR TF(DELL)
      DELL=VA-LA
      CUL(J+1)=XK*DELL*SQR TF(DELL)
      J=J+2
      K=K+JD
      IF(J-NAJ) 39,4C,4C
      END

```

```

SUBROUTINE TRAJY(KE,KED)
C TRAJY CALCULATES THE COORDINATES IN THE Y SWEEP
COMMON URH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1  NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DX,DELY,YEF,XCM,F,KISW,ETX,
2  ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCR,KCH,HSL,XSL,EPS,NPOL,
3  NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUP,
4  RFDOWN,XCU,NOT,NO,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5  IXO,HGH,XLOH,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1  RH(4CCC),L(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2  ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3  ETX(4C),ETY(4C),XCU(4C),KCY(10)
JE=KE
JED=KED
JD=NPONT
CO 40 K=1,NTJ
41 AD=AX(K)/H
JX=AD
XA=JX
XA=AD-XA
JP=JE+JX*JD
JS=4
JQ=JP-1
JOP=JQ+JD
UL=(1.-XA)*L(JQ)+XA*L(JOP)
JAP=JP+JD
UK=(1.-XA)*L(JP)+XA*L(JAP)
IF(XA-.5) 1,1,5
5 XA=XA-1.
JX=JX+1
JQ=JQ+13
1 JOP=JQ+JD
JAP=JQ-JD
YLA=(XA+.5)*L(JCP)-2.*XA*L(JQ)+(XA-.5)*U(JAP)
DY=VX(K)/VY(K)
JDX=JX
7 XB=XA+DY
IF(DY) 5,9,10
8 XB=1.+XB
JX=JX-1
9 IF(XB+.5) 8,12,12
11 XB=XB-1.
JX=JX+1
10 IF(XB-.5) 12,12,11
12 JP=JE+JX*JD
JOP=JP+JD
JAP=JP-JD
YRA=(XB+.5)*U(JCP)-2.*XB*L(JP)+(XB-.5)*U(JAP)
IF(XB) 19,20,20
19 XB=1.-ABSF(XB)
JP=JP-13
20 JQ=JP-1
JOP=JQ+JD
JAP=JP+JD
LSN=(1.-XB)*U(JQ)+XB*L(JOP)
LQN=(1.-XB)*U(JP)+XB*L(JAP)

```

```
DUX=.5*(UK-LL-LSN+LQN)
VYB=SQR TF(VY(K)**2-2.*XQM*DUX)
DT=2.*DX/(VYB+VY(K))
JS=JS-1
YA=.5*XQM*(YLA+YRA)/DX
DY=DT*(VX(K)-.5*YA*DT)/DX
JX=JUX
IF(JS) 21,21,7
21 VY(K)=VYB
VX(K)=VX(K)-YA*DT
AX(K)=AX(K)+DY*DX
40 CONTINUE
RETURN
END
```

```

SUBROUTINE CALRY(KE,KED)
C CALRY CALCULATES THE RHS IN THE Y SWEEP
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREG,NTP,
1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,CX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCLCR,KCH,HSL,XLSL,EPS,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBA,RHUP,
4 RHDOWN,XCL,NOT,NC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5 IXO,FGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(1C),URH(4000),UB(4000),JT(4000),
1 RH(4000),U(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCL(4C),KCY(10)
JE=KE
JED=KED
DO 33 JJ=1,NA J
J=JJ
1 HT=AX(J+1)-AX(J)
IF(HT) 12,12,13
12 HT=-HT
XH=AX(J+1)
JX=XH/DELY
NB=AX(J)/DELY+1.
WA=SQR TF(VX(J+1)**2+VY(J+1)**2)
WB=SQR TF(VY(J)**2+VY(J)**2)
CO TO 14
13 XH=AX(J)
JX=XH/DELY
WA=SQR TF(VX(J)**2+VY(J)**2)
WB=SQR TF(VX(J+1)**2+VY(J+1)**2)
NB=AX(J+1)/DELY+1.
14 YL=XH+HT
DO 32 KK=JX,NB
K=JE+KK*NPOINT
XA=KK
XLD=(XA-.5)*DELY
XX=XLD+DELY
YU=MAX1F(XLD,XH)
YD=MIN1F(XX,YL)
IF(YD-YL) 32,32,27
27 XA=.5*(YD+YL)
W=WA+(XA-XH)*(WB-WA)/HT
RH(K)=RH(K)+(YD-YL)*CL(JJ)/(HT*W)
32 CONTINUE
33 CONTINUE
RETURN
END

```

```

SUBROUTINE TRAJX(KE,KED,BX)
C TRAJX PICKS UP THE TRAJECTORIES IN THE X SWEEP
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1 NPIT,XEP,NXEP,CL,NPENT,AX,AY,VX,VY,DY,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCR,KCH,HSL,XSL,EPS,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUP,
4 RFDOWN,XCL,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5 IXO,HGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1 RH(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40),
2 ,VX(40),KCH(40),PTX(40),PTY(40),AY(40),AX(40),ATX(40),ATY(40),
3 ETX(40),ETY(40),XCL(40),KCY(10)

JE=KE
JED=KED
AAX=BX
DO 10 J=1,NTJ
10 IF(AX(J)+1.) 10,10,1
1 IF(KCH(J)) 2,11,11
2 IF(AAX-AX(J)) 10,10,3
3 IF(J-1) 4,4,5
5 CU(J-1) =-CL(J-1)
4 JX=(AY(J)+.CC1)/DX
IF(NOT) 40,40,41
40 IF(JOT) 41,41,42
42 NOT=JOT
41 JP=JE+JX
JQ=JP-NPONT
JS=4
XA=AAX-AX(J)
UL=(1.-XA)*L(JP)+XA*L(JQ)
UK=U(JP)
UA=(1.-XA)*L(JP-1)+XA*L(JQ-1)
UD=(1.-XA)*L(JP+1)+XA*U(JQ+1)
YLA=.5*LD-.5*LA
CY=VY(J)/VX(J)
JDX=JX
7 XB=0Y
GO TO 30
31 XB=XB-1.
JX=JX+1
30 IF(XB-.5) 12,12,31
12 JP=JE+JX
YRA=(XB+.5)*L(JP+1)-2.*XB*U(JP)+(XB-.5)*U(JP-1)
IF(XB) 19,19,20
19 JP=JP-1
XB=1.-ABS(XB)
20 JQ=JP-NPONT
UQN=(1.-XB)*U(JP)+XB*L(JP+1)
USN=(1.-XB)*U(JQ)+XB*L(JQ+1)
LSN=(1.-XA)*UQN+XA*LSN
DLX=.5*(UK-LL+UQN-USN)
VXB=SQR(TF(VX(J)**2-2.*XQM*DLX))
DT=2.*XA/(VXB+VX(J))
JS=JS-1
YA=.5*XQM/DX*(YLA+YRA)
CY=DT*(VY(J)-.5*YA*DT)/DX

```

```
JX=JOX
IF(JS) 21,21,7
21  VX(J)=VXB
    VY(J)=VY(J)-YA*DT
    AY(J)=AY(J)+DY*DX
    AX(J)=AAX
    KCH(J)=C.
    GO TO 1C
11  AX(J)=AX(J)+DX
C TRAJ CALCULATES THE TRAJECTORY COORDINATES IN THE X SWEEP
    CALL TRAJ(JE,JED,J)
1C  CONTINUE
    RETURN
    END
```

```

SUBROUTINE TRAJ(KE,KED,M)
C IRAJ CALCULATES THE TRAJECTORY COORDINATES IN THE X SWEEP
COMMON LRH,LB,JI,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREG,NTP,
1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCR,KCH,HSL,XLSL,EPS,NPOL,
3 NPUL,JOI,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RFUP,
4 RFDOWN,XCU,NOT,MO,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJOT,
5 IXO,HGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),LRH(4000),UB(4000),JT(4000),
1 RF(4000),U(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(4C),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCL(4C),KCY(10)
JE=KE
JED=KED
K=M
C TRAJ CALCULATES COORDINATES FOR X-SWEEP
AD=AY(K)/DX
JX=AD
XA=JX
XA=AD-XA
JP=JX+JE
JS=4
JQ=JP-NPONT
UL=(1.-XA)*L(JQ)+XA*L(JQ+1)
UK=(1.-XA)*L(JP)+XA*L(JP+1)
IF(XA-.5) 1,1,5
5 XA=XA-1.
JQ=JQ+1
JX=JX+1
1 YLA=(XA+.5)*U(JQ+1)-2.*XA*U(JQ)+(XA-.5)*U(JC-1)
CY=VY(K)/VX(K)
JOX=JX
7 XB=XA+DY
IF(CY) 5,9,10
8 XB=XB+1.
JX=JX-1
9 IF(XB+.5) 8,12,12
11 XB=XB-1.
JX=JX+1
10 IF(XB-.5) 12,12,11
12 JP=JE+JX
YRA=(XB+.5)*U(JP+1)-2.*XB*U(JP)+(XB-.5)*U(JP-1)
IF(XB) 19,20,20
19 JP=JP-1
XB=1.-ABSF(XB)
20 JQ=JP-NPONT
USN=(1.-XB)*U(JQ)+XB*L(JQ+1)
UQN=(1.-XB)*U(JP)+XB*L(JP+1)
DX=.5*(UK-UL-USN+UCN)
VXB=SQR TF(VX(K)**2-2.*XQM*DLX)
DT=2.*DX/(VXB+VX(K))
JS=JS-1
YA=.5*XQM*(YLA+YRA)/DX
DY=DT*(VY(K)-.5*YA*DT)/DX
JX=JOX
IF(JS) 21,21,7

```

```
21      VX(K)=VXB  
      VY(K)=VY(K)-YA*DT  
      AY(K)=AY(K)+DY*DX  
      RETURN  
      END
```

```

SUBROUTINE CALRX(KE,KED,KD)
C CALRX CALCULATES THE RHS IN THE X SWEEP
COMMON LRH,LB,JI,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1 NPIT,XEP,NXEP,CL,APCNT,AX,AY,VX,VY,LX,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EPS,NPCL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUP,
4 RHDOWN,XCL,NUT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5 IXO,FGH,XLOW,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),LRH(4000),UB(4000),JT(4000),
1 RF(4CC0),L(4CCC),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCL(4C),KCY(10)

JAD=KD
JE=KE
JED=KED
DO 33 JJ=1,NAJ
J=JJ
IF(CL(JJ)) 33,33,1
1 FT=AY(J+1)-AY(J)
IF(FT) 12,12,13
12 FT=-HT
XH=AY(J+1)
JX=XH/DELY
NB=AY(J)/DELY+1.
WA=SQR TF(VX(J+1)**2+VY(J+1)**2)
WB=SQR TF(VX(J)**2+VY(J)**2)
CU TO 14
13 XF=AY(J)
JX=XH/DELY
WA=SQR TF(VX(J)**2+VY(J)**2)
WB=SQR TF(VX(J+1)**2+VY(J+1)**2)
NB=AY(J+1)/DELY+1.
14 NA=JX+JE
NB=NB+JE
YL=XH+HT
DO 32 K=NA,NB
IF(K-JAD) 32,32,1C
1C XA=K-JE
XUD=(XA-.5)*DELY
XX=XUD+DELY
YL=MAX1F(XUD,XH)
YD=MIN1F(XX,YL)
IF(YD-YL) 32,32,27
27 XA=.5*(YD+YL)
W=WA+(XA-XH)*(WB-WA)/HT
RF(K)=RF(K)+(YD-YL)*CL(JJ)/(HT*W)
32 CONTINUE
33 CONTINUE
RETURN
END

```

```

SUBROUTINE CORRCT(BAX,M)
C CORRCT CONDITIONALLY ENDS OR SHIFTS THE TRAJECTORIES
COMMON LRH,LB,J1,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,KREG,NTP,
1 NPIT,XEP,NXEP,CL,APCNT,AX,AY,VX,VY,DY,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCR,KCH,HSL,XLSL,EPN,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KKL,KAN,LC,LD,VAT,VBT,SIZE,KBN,RHUP,
4 RFDOWN,XCL,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABB,NURL,NJCT,
5 IXO,HGH,XLOW,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1 RH(4000),U(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(40)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCU(4C),KCY(10)
      JN=M
      AAX=BAX
      IF(JN-12) 1C,1,12
12     IF(JN-13) 1C,13,1C
1      CU 30 J=1,NTJ
      IF(AX(J)+1.) 30,3C,3
3      IF( XLOk-AY(J)) 4,3C,30
4      AYD=AY(J)-AY(J+1)
      AY(J)= XLOW
      IF(AY(J)-AY(J+1)) 5,5,6
5      AY(J)=-1.
      AX(J)=-1.
      CU(J)=0.
      GO TO 30
6      AYDS=AY(J)-AY(J+1)
      AYDL=AYD-AYDS
      CL(J)=CL(J)*AYDS/AYD
      VY(J)=(VY(J)*AYDS+VY(J+1)*AYDL)/AYD
      VX(J)=(VX(J)*AYDS+VX(J+1)*AYDL)/AYD
30     CONTINUE
10     RETURN
13     CU 40 J=1,NTJ
      IF(AX(J)+1.) 40,4C,21
21     IF(AY(J)- HGH ) 22,4C,4C
22     AYD=AY(J-1)-AY(J)
      AY(J)= HGH
      IF(AY(J)-AY(J-1)) 24,23,23
23     AY(J)=-1.
      AX(J)=-1.
      CU(J-1)=0.
      GO TO 4C
24     AYDS=AY(J-1)-AY(J)
      AYDL=AYD-AYDS
      CU(J-1)=CU(J-1)*AYDS/AYD
      VY(J)=(VY(J)*AYDS+VY(J-1)*AYDL)/AYD
      VX(J)=(VX(J)*AYDS+VX(J-1)*AYDL)/AYD
4C     CONTINUE
      GO TO 1C
      END

```

```

SUBROUTINE RTEST
COMMON LRH,LB,JT,RH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1 NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DXY,DELY,YEP,XCM,F,KISW,ETX,
2 ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,ACCCR,KCH,HSL,XLSL,EPSP,NPOL,
3 NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VET,SIZE,KBN,RHUP,
4 RHDOWN,XCL,NOT,MC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5 IXO,HGH,XLOH,XMPR,NEM
DIMENSION KBN(2C),LD(4),LC(10),URH(4000),UB(4000),JT(4000),
1 RF(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2 ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3 ETX(4C),ETY(4C),XCU(4C),KCY(10)

C RTEST CHECKS ON THE UPPER AND LOWER BOUNDS
C KAN=TEST POINT
  KAN=KAN
  RT=RH(KAN)
  SX=1.

C FIRST TWO CYCLES SET UPPER AND LOWER BOUNDS
  IF(KRL) 22,23,22
  23  IF(MO) 11,22,11
  22  IF(KRL-NRL+1) 5,2,1

C RHUP=UPPER BOUND
  1  RHUP=RT
    GO TO 19
  2  IF(RT-RHUP) 18,18,3

C RHDOWN=LOWER BOUND
  3  RHDOWN=RHUP
    MO=-MO
    RHUP=RT
  4  GO TO 19

C MO=+1 ,CHECK IS ON UPPER BOUND
C MO=-1 ,CHECK IS ON LOWER BOUND
C MO=0 ,NO CHECK
  5  IF(MO) 10,19, 8
  6  SX=SX*RX
    DO 7 J=1,NTOP
    7  RH(J)=RH(J)*RX
    8  IF(RHUP-RH(KAN)) 6,9,9
    9  IF(RH(KAN)-RHDOWN) 11,16,16
   10  IF((RT-RHDOWN)*(RT-RHUP)) 16,16,11
   11  POW=1./SX
    DO 12 J=1,NTOP
   12  RH(J)=.5*(RH(J)*POW+LRH(J))
    KRL=1
    JJ=NKL+1
    KCY(JJ)=777
    KCY(JJ+1)=777
    WRITE OUTPUT TAPE 6,100
    IF(MO) 14,13,13
   13  RHUP=RT
    GO TO 15
   14  RHDOWN=RT
   15  MO=0
   16  IF(MO) 18,19,17
   17  RHUP=RH(KAN)
    GO TO 19
   18  RHDOWN=RT

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```
19      XT(1)=    RHDOWN
          XT(2)=    RH(KAN)
          XT(3)=    RHLR
          WRITE OLTPUT TAPE 6,1C1,(XT(J),J=1,3), U(KAN)
          IF(IXO) 21,21,24
24      IW=NRL-KRL+1
          IF(KCY(IW)) 21,21,2C
C TROUT PRINTS OUT THE RHS
20      CALL TROUT
21      MO=-MO
          KRL=KRL-1
          RETURN
1C0      FORMAT(11H0RH=AVERAGE    )
1C1      FORMAT(7HCRHLOW=F6.3,4H RH=F6.3,6H RHUP=F6.3,7H UTEST=F11.4)
          END
```

```

      SUBROUTINE TROUT
C TROUT PRINTS OUT THE RHS
      COMMON LRH,LB,JT,KH,L,XT,KT,LINC,LB,XR,NTCP,NUL,NLIN,NREC,NTP,
1      NPIT,XEP,NXEP,CL,NPCNT,AX,AY,VX,VY,DX,DELY,YEP,XCM,F,KISW,ETX,
2      ETY,PTY,PTX,NAJ,NTJ,KBA,VA,VB,VC,NCCCR,KCH,HSL,XLSL,EFS,NPOL,
3      NPUL,JOT,KAB,NSPAN,RX,NRL,KRL,KAN,LC,LD,VAT,VBT,SIZE,KEN,RHUP,
4      RHDOWN,XCU,NOT,NC,KCY,NSWP,ATX,ATY,KAT,KATT,KABE,NURL,NJCT,
5      IXO,FGH,XLOW,XMPR,NEM
      DIMENSION KBN(2C),LD(4),LC(10),LRH(4000),UB(4000),JT(4000),
1      RH(4000),L(4000),XT(630),KT(630),LINC(50),LB(200),CU(40),VY(4C)
2      ,VX(4C),KCH(4C),PTX(4C),PTY(40),AY(40),AX(40),ATX(40),ATY(4C),
3      ETX(4C),ETY(4C),XCL(4C),KCY(10)
      IF(MO) 10,11,10
10     K=1
      CO 1 J=1,NTOP
      KT(K)=J
      XT(K)=RH(J)
      K=K+1
      IF(K-S) 1,2,2
2      WRITE OLTPUT TAPE 6,100,(KT(L),L=1,8),(XT(L),L=1,8)
      K=1
1      CONTINUE
      IF(K-2) 8,6,6
6      CO 7 J=K,8
      KT(J)=C
7      XT(J)=C.
      WRITE OLTPUT TAPE 6,100,(KT(L),L=1,8),(XT(L),L=1,8)
8      RETURN
11     WRITE OLTPUT TAPE 6,1C1
      K=1
      CO 21 J=1,NTOP
      KT(K)=J
      XT(K)=.5*(RH(J)+LRH(J))
      K=K+1
      IF(K-S) 21,22,22
22     WRITE OLTPLT TAPE 6,100,(KT(L),L=1,8),(XT(L),L=1,8)
      K=1
1      CONTINUE
      IF(K-2) 8,26,26
26     CO 27 J=K,8
      KT(J)=C
27     XT(J)=C.
      WRITE OLTPLT TAPE 6,100,(KT(L),L=1,8),(XT(L),L=1,8)
      GO TO 8
1C0    FORMAT(1H 8I5,8F1C.3)
1C1    FORMAT(21H ORH PRINT IS AVERAGE.    )
      END

```

\* DATA

11

1 HAMZA/AXI-SYMETRIC ION ENGINE

MAXIMUM SLOPE TEST IS ON 1.

MESH SIZE=.005

EMITTER POTENTIAL =5657.2 VOLTS

UPPER ELECTRODE=-3184.8 VOLTS

LOWER ELECTRODE=0. VOLTS

SUPPRESSION FACTOR =.4

CURRENT IS CALCULATED BY EQUAL DELTA X

INITIAL SWEEP IS IN Y-DIRECTION

EMITTER COORD. ARE TAKEN AS THE BEGINNING TRAJ. CCCRD.

1

221	13	106	1	30	200	1									
50	20	60	308	221	0										
0	1	-1	1	-1	1	-1	1	-1	1	-1	1	-1	1	-1	-1
1	1	-1	-1	1	0										
0	1	13	0	13	5	10	0	26	4	11	0	39	3		
11	0	52	2	12	0	65	2	12	0	78	2	12	0		
91	2	12	0	104	2	12	0	117	2	12	0	130	2		
12	0	143	2	12	0	156	2	10	0	156	12	12	0		
169	3	6	0	169	9	13	0	182	3	5	0	182	6		
13	0	195	6	13	0	208	1	13	0						
0	-13	13	-1	0	1										
.95		.95		.95		-.975		3.8		-.925			9		
0	-13	13	-1	0	1										
.90		.90		.90		-.925		3.6		-.875			15		
0	-13	13	-1	0	1										
.85		.85		.85		-.875		3.4		-.825			21		
0	-13	13	-1	0	1										
.80		.80		.80		-.825		3.2		-.775			27		
0	-13	13	-1	0	1										
.75		.75		.75		-.775		3.0		-.725			33		
0	-13	13	-1	0	1										
.70		.70		.70		-.725		2.8		-.675			39		
0	-13	13	-1	0	1										
.65		.65		.65		-.675		2.6		-.625			45		
0	-13	13	-1	0	1										
.60		.60		.60		-.625		2.4		-.575			51		
0	-13	13	-1	0	1										
.55		.55		.55		-.575		2.2		-.525			57		
0	-13	13	-1	0	1										
.50		.50		.50		-.525		2.0		-.475			63		
0	-13	13	-1	0	1										
.45		.45		.45		-.475		1.8		-.425			69		
0	-13	13	-1	0	1										
.40		.40		.40		-.425		1.6		-.375			75		
0	-13	13	-1	0	1										
.49375		.49375		.49375		0.		1.9625		-.975			81		
0	-13	13	-1	0	1										
.20625		.20625		.20625		-.425		.8375		0.			87		
0	-13	13	-1	0	1										
.65		.65		.65		.675		2.6		-.625			93		
0	-13	13	-1	0	1										
.75		.75		.75		-.775		3.0		.725			99		
0	-13	13	-1	0	1										

.75084622	1.9729456	.55555975	.93708046	4.6810257	-.775	105
0 -13	13 -1	C 1				
.75	.95096646	.81002600	-.775	3.2609924	-.725	111
0 -13	13 -1	C 1				
.60	.90187053	.68290213	-.625	2.7847726	-.575	117
0 -13	13 -1	C 1				
.38358187	2.2193357	.61003368	-.575	4.7723274	1.3679580	123
0 -13	13 -1	C 1				
.33791599	1.1218393	.48075906	-.525	3.5244100	1.3968117	129
0 -13	13 -1	C 1				
.85963256	.85856154	.85856154	1.0136909	3.6058140	-.875	135
0 -13	13 -1	C 1				
.72163908	.89234030	.76989569	1.8545964	4.4418322	-.925	141
0 -13	13 -1	C 1				
.30196656	.30748522	.30748522	-.475	2.3804767	1.2905064	147
0 -13	13 -1	C 1				
.36033430	.36432314	.36432314	-.475	1.9267936	.72314733	153
0 -13	13 -1	C 1				
.40067118	.40311125	.40311125	-.475	1.8329834	.55136090	159
0 -13	13 -1	C 1				
.42863746	.42976792	.42976792	-.475	1.8067430	.47220713	165
0 -13	13 -1	C 1				
.78097865	.77732208	.77732208	1.4997761	3.9794202	-.925	171
0 -13	13 -1	C 1				
.59066353	.58478411	.58478411	3.9264021	6.0209701	-.925	177
0 -13	13 -1	C 1				
.52532034	.67746792	1.2721693	8.9925431	11.867200	-.925	183
0 -13	13 -1	C 1				
.86194494	.86093254	.86093254	1.0081411	3.6050062	-.875	189
0 -13	13 -1	C 1				
.79513451	.79256099	.79256099	1.1984555	3.6585774	-.875	195
0 -13	13 -1	C 1				
.36875	.36875	.36875	C.	1.4625	-.725	201
0 -13	13 -1	C 1				
.70	.70861765	.72628258	-.725	2.8349002	-.675	207
0 -13	13 -1	C 1				
.65	.73865090	.57195333	-.675	3.0106042	-.625	213
0 -13	13 -1	C 1				
.70771769	.88514481	1.4700828	-.775	3.9505388	.82031120	219
0 -13	13 -1	C 1				
.45	.45	.45	.475	1.8	.425	225
0 -13	-13 -1	C 1				
.85	.85	.85	-.875	3.4	-.825	231
0 -13	13 -1	C 1				
.50	1.7845933	.71418604	-.525	3.4987793	-.475	237
0 -13	-11 -1	C -11				
.60710078	.85717620	2.00065597	-.825	5.2088771	1.5201413	243
0 -12	13 -12	C 1				
.44784299	4.4455398	.75533228	3.8860436	9.8123157	-.7250000	249
0 -13	13 -1	C 1				
.70000000	.97871180	.77937012	-.7250000	3.1580819	-.6750000	255
0 -13	13 -1	C 1				
.65	.46800967	.58034818	-.675	2.3483578	-.625	261
0 -13	13 -1	C 1				
.85	.85	.85	.875	3.4	-.825	267
0 -13	-13 -1	C 1				

.95	.95	.95	.95	.975	3.8	-.925	273
0	-13	13	-1	0	1		
.45	.45	.45	.45	-.475	1.8	.425	279
0	-13	13	-1	0	1		
.20625	.20625	.20625	.20625	-.425	.8375	0.	285
0	-13	13	-1	0	1		
.55	.55	.55	.55	-.575	2.2	.525	291
0	-13	13	-1	0	1		
.376800000	.370959599	.370959599	2.3687500	3.6856691	-.5750000		297
0	-13	13	-1	0	1		
.50	.50	.50	.50	-.525	2.0	.475	303
0	0	0	0	0	0	0	0
0	0	0	105	111	39	45	117
0	0	0	267	27	33	39	45
0	0	135	21	27	33	39	45
0	141	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	273	15	21	27	33	39	45
0	171	15	21	27	33	39	45
0	171	15	21	27	33	39	45
0	183	15	21	27	33	207	213
0	0	189	21	27	219	0	297
0	0	195	231	243	249	255	45
0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0
.93417927							
1	1	13	1				
5657.2							
1	1	13	1				
5657.2							
1	1	13	1				
5657.2							
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
3	1	4	5				
5657.2	0.			5657.2			
2	1	6	6				
5657.2	0.						
1	12	13	1				

5657.2  
 2 1 4 4  
 5657.2 -3184.8  
 1 9 13 1  
 0.  
 1 1 5 1  
 5657.2  
 1 6 13 1  
 -3184.8  
 5 1 4 4  
 5657.2 0. 0. 0.  
 1 1 13 1  
 0.  
 1 119 9 9 67 13 11 2 200 7 1 5 7 4  
 6 80 93 106 119 132 67  
 1 11 15 2 14 210 13  
 14 13 16 1  
 5657.2 -3184.8 0. .005 .0001 .40  
 1. 0. .03872 .04777 8. 6.  
 8.854E-12 9.649E+07 132.91E+0  
 6000. -3000. 500.  
 .0250 .0275 .0300 .0325 .0350 .0375 .0400  
 .0425 .0450  
 0. 0.  
 0.  
 1

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